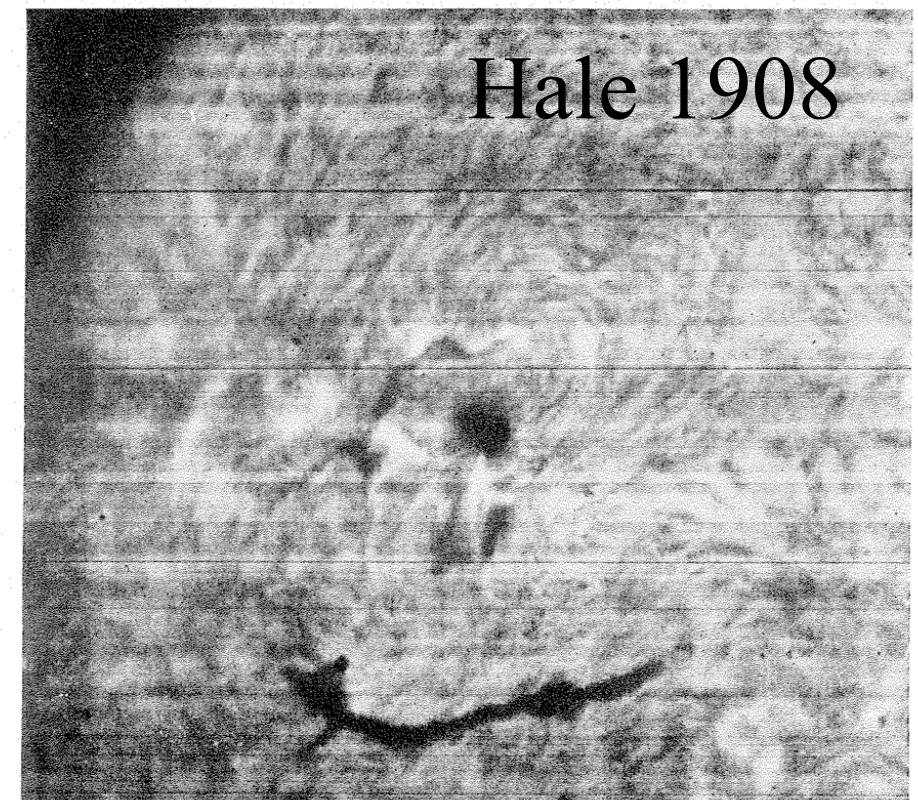


# Physics of the chromosphere and the lower coronal boundary conditions

Philip Judge

HAO, NCAR, Boulder CO USA

PLATE X



December 2008

FIG. 1.—SUN-SPOT AND HYDROGEN ( $H\alpha$ ) FLOCCULI  
1908, May 29, 4<sup>h</sup> 26<sup>m</sup> P. M. Scale: Sun's Diameter = 0.3 Meter



NCAR

The National Center for Atmospheric Research is operated by the University Corporation for Atmospheric Research under sponsorship of the National Science Foundation. An Equal Opportunity/Affirmative Action Employer.



# the chromosphere:

## interface between photosphere and corona

- partially ionized: thermostat
- stratified: spans 9 pressure scale heights
- so it usually contains plasma  $\beta=1$  surface
  - forced at the base
  - force-free at the top
- requires 30-100x as much power as the corona
- **is the lower boundary for the corona**
  - modulates flow of mass, momentum, energy and magnetic field into the corona
  - implicit mass reservoir in coronal loop scaling laws



## Gold (1964)

- chromosphere occupies thick black line
- the electro-dynamics of the **chromosphere** is **critical** to the supply of magnetic free energy into the corona.
- traditionally it is treated as in the figure

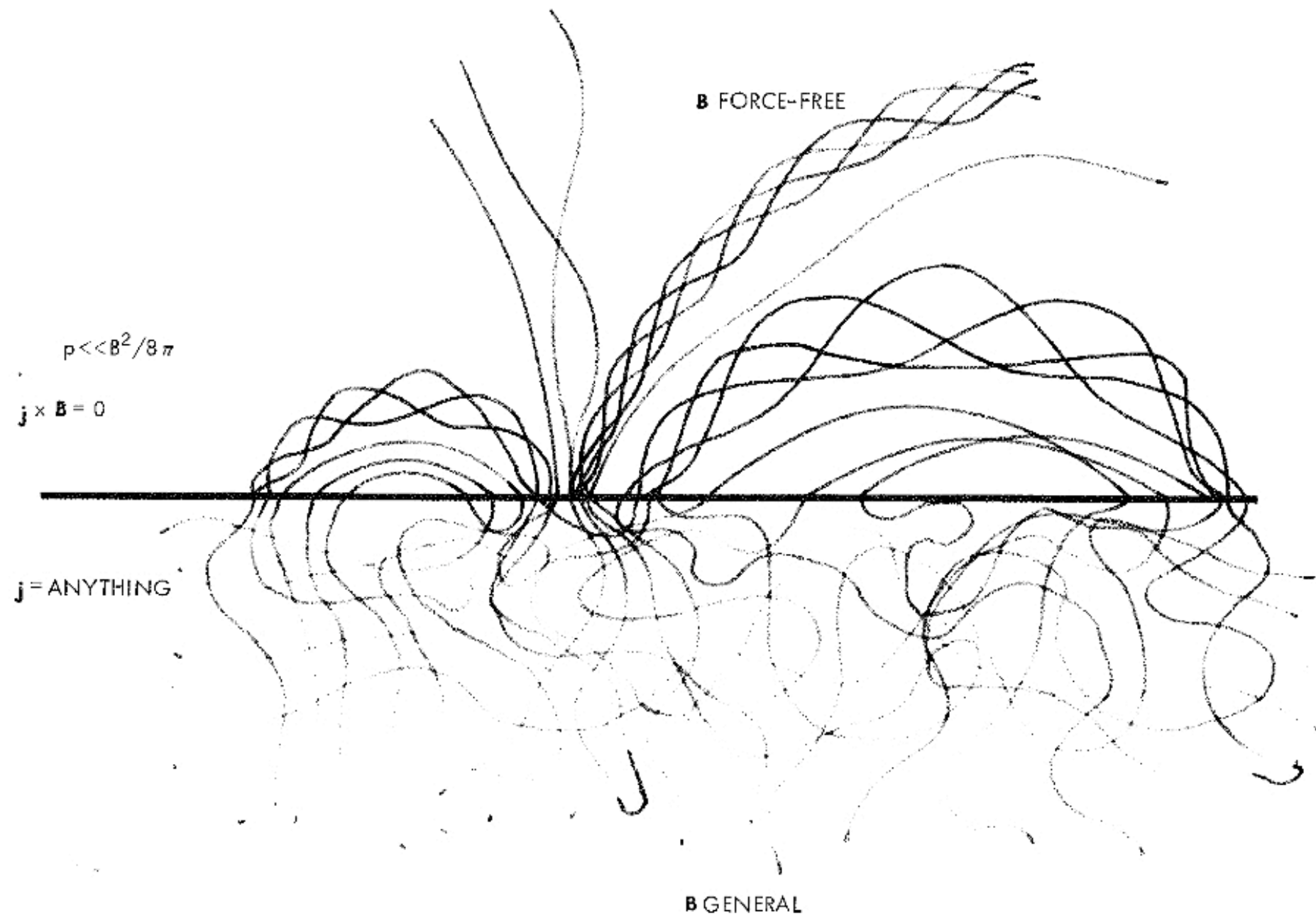
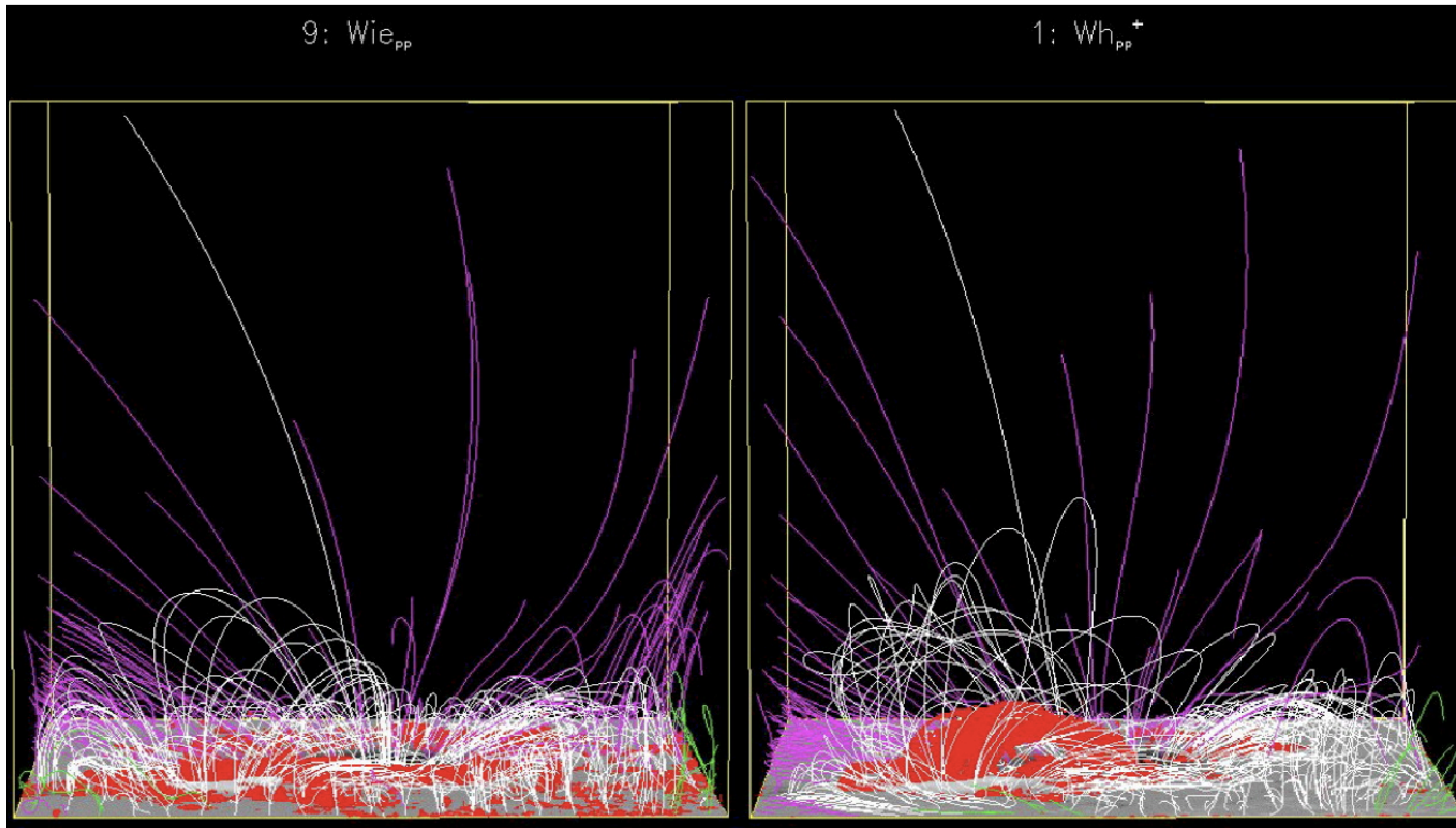


FIGURE 44-2. Magnetic field in a turbulent conducting medium. The fluid pressure is assumed large compared with magnetic forces below the dividing plane and small above it.

# nlff field extrapolation (Schrijver et al 2008)

red:  
current



force free extrapolations from photospheric vector polarimetry. Photospheric boundary is *not* force-free...

# DOT and TRACE 9 Jul 2005 (A.G. de Wijn, R. J. Rutten)

photosphere  
- forced

top of  
chromosphere,  
corona  
- force free



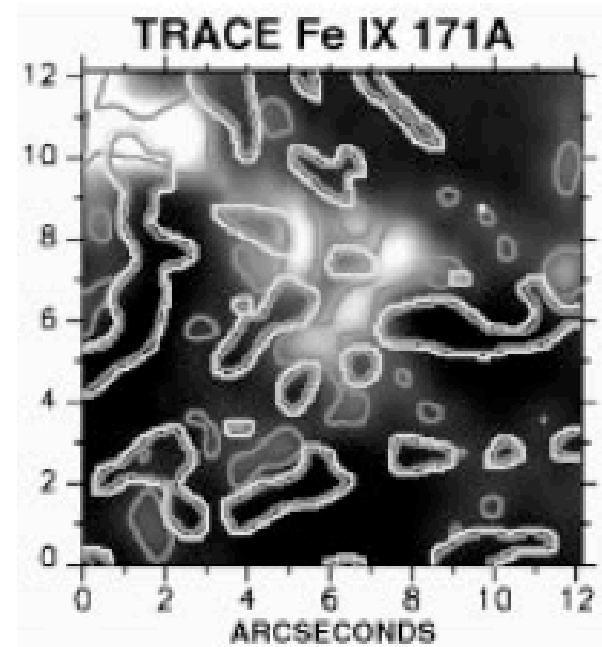


# failure to trace magnetic field lines through the solar atmosphere

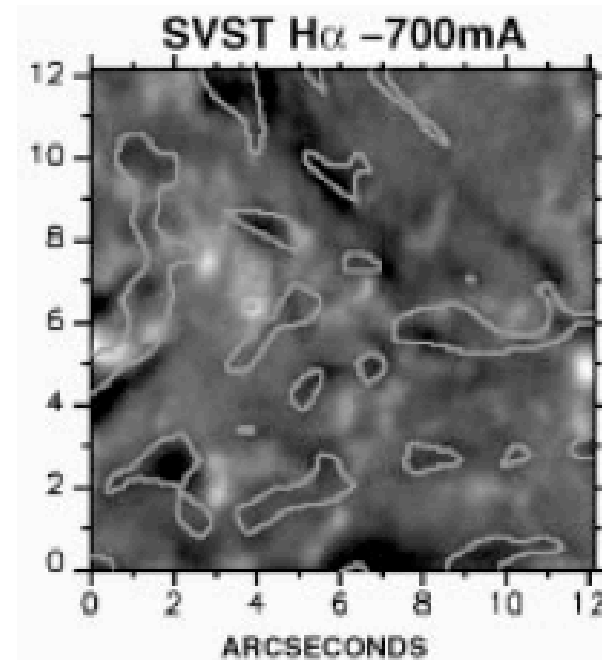
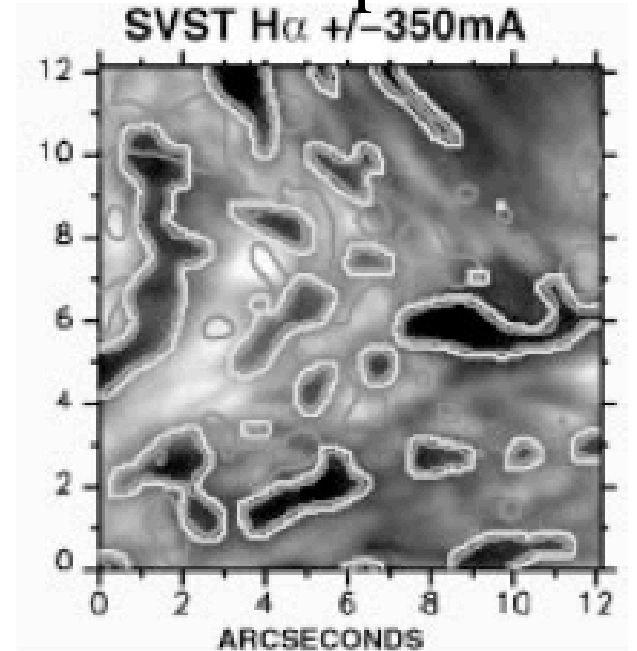
- high plasma conductivity allows tracing of field lines using
  - photospheric observations (polarimetry, proxies)
  - coronal threads
- **TRACE** & other missions failed to do this
- **why?- chromosphere**

De Pontieu et al. 1999  
“moss”

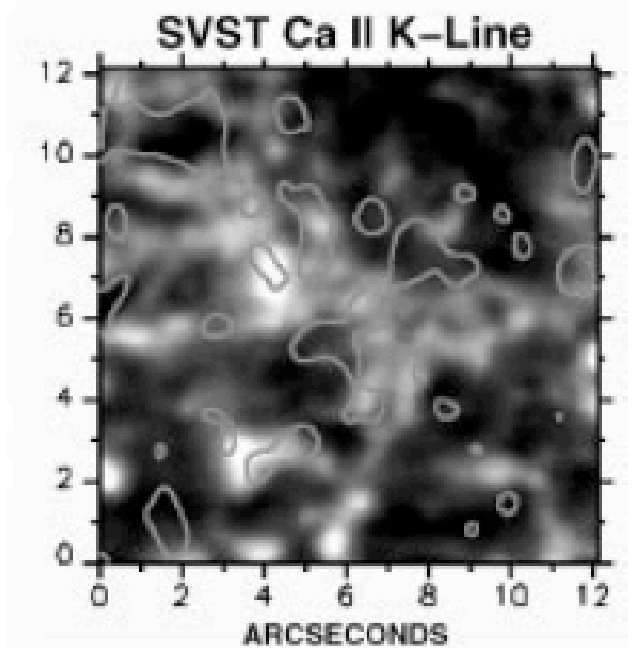
lower corona



upper  
chromosphere



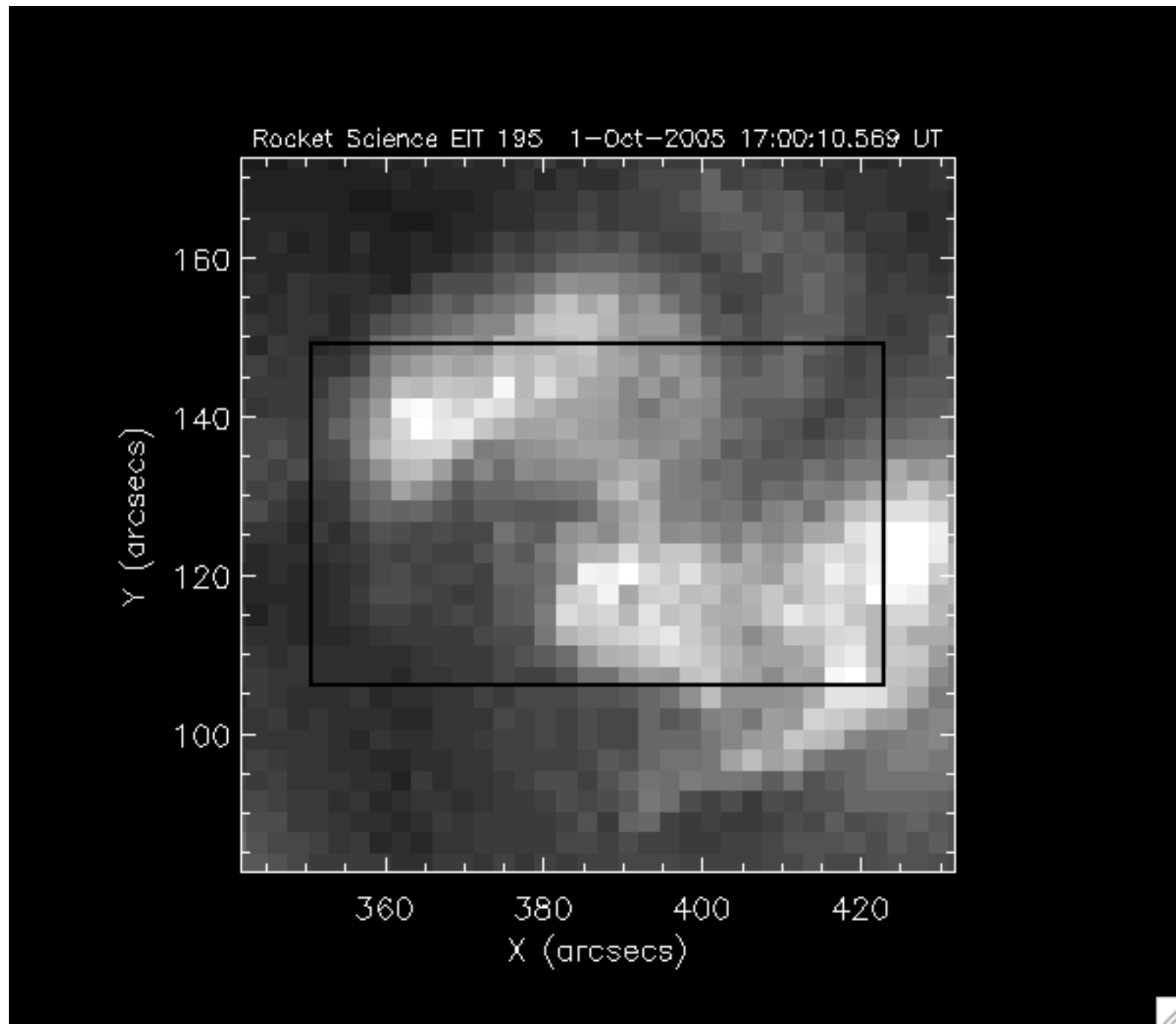
upper  
chromosphere



magnetic elements +  
reverse granulation

**magnetic interface  
observations:  
an example**

# Small AR, pores

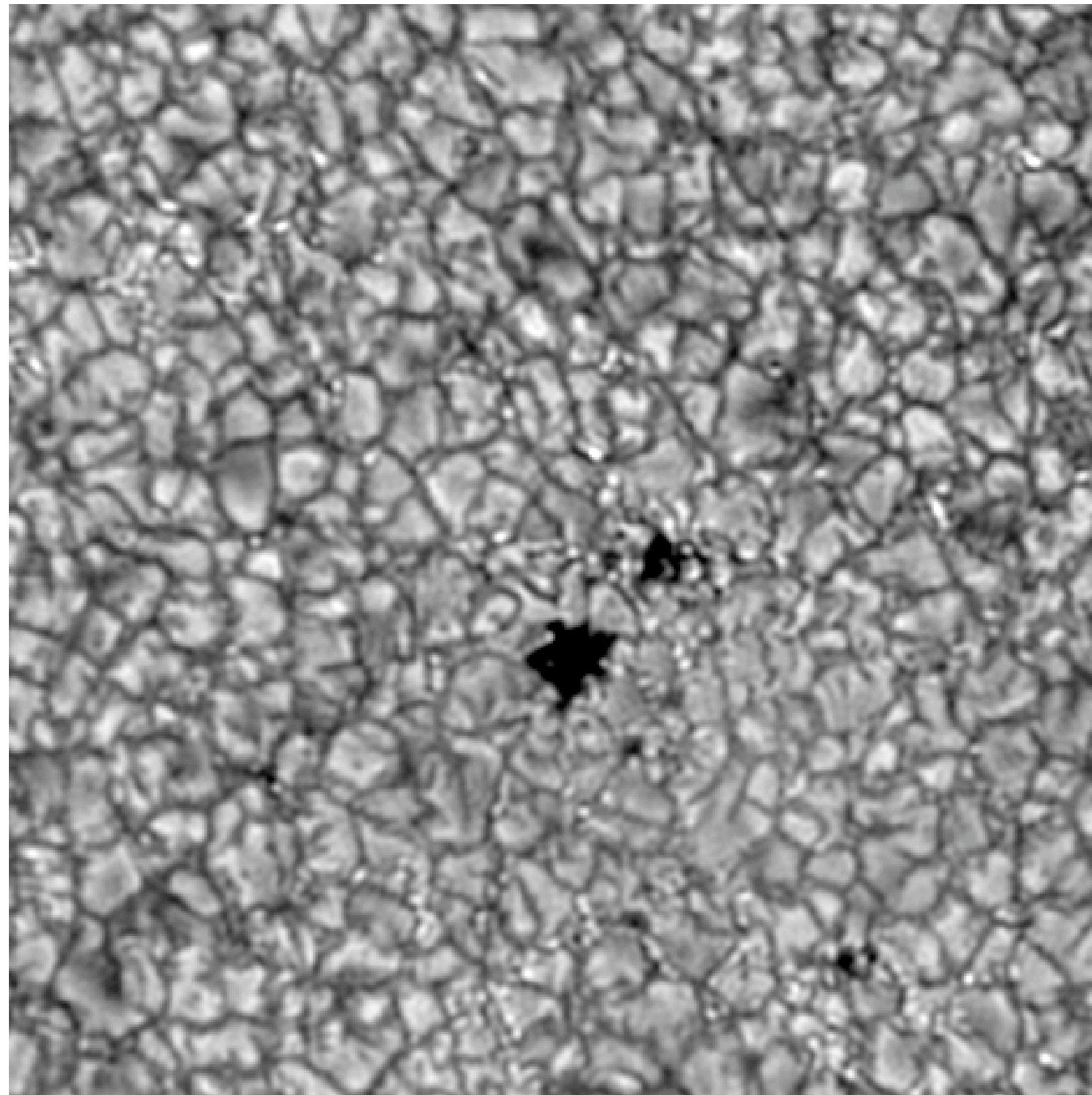




# Chromosphere as seen with IBIS

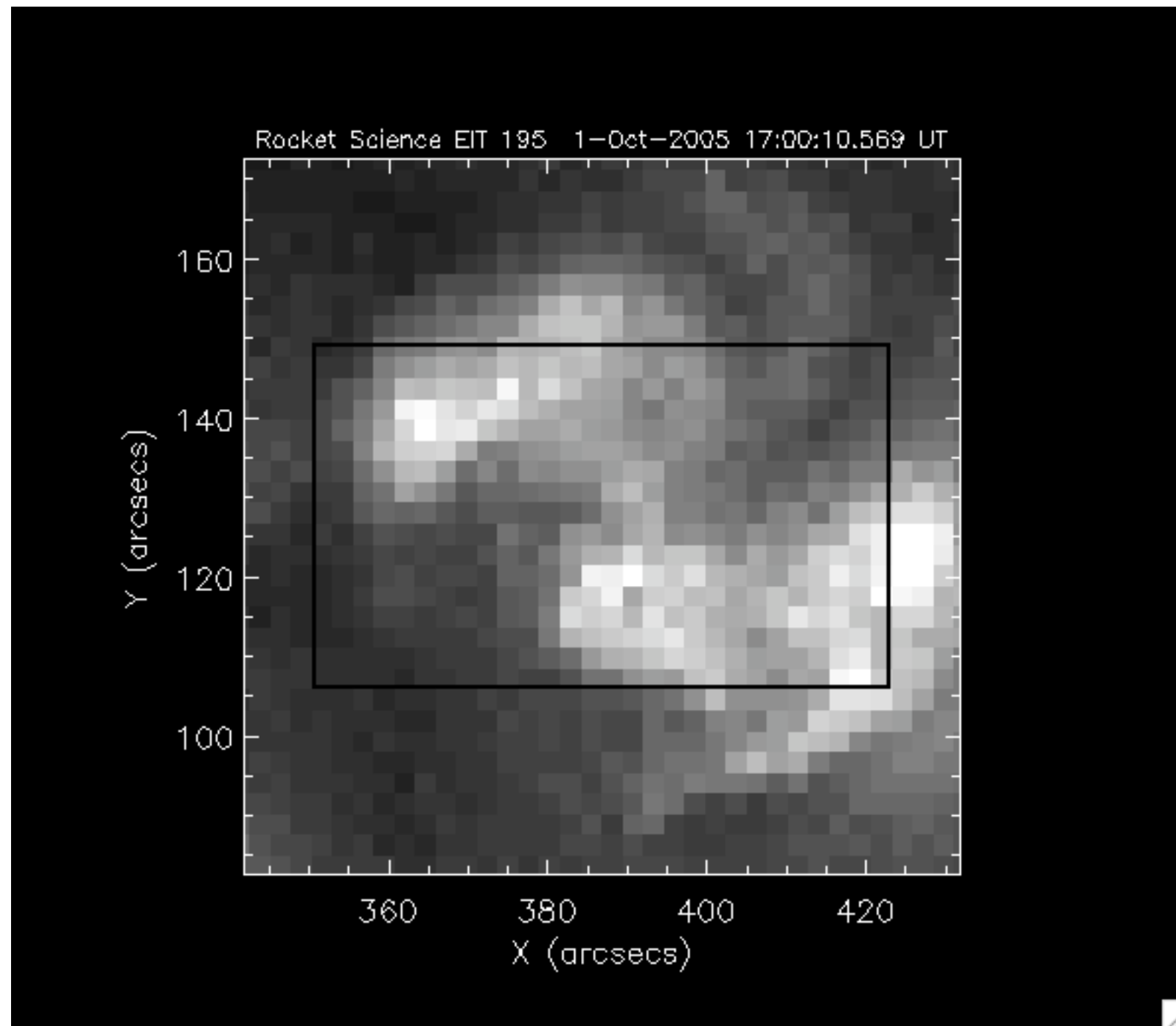
- Ca II 854.2 nm
- samples many pressure scale heights
- high resolution
  - resolution  $\approx 0.3''$  (DST limit  $0.24''$ )
  - (FOV  $40'' \times 40''$ )

base of corona is  
**very** different from  
photosphere



G. Cauzzi et al 2008, A+A

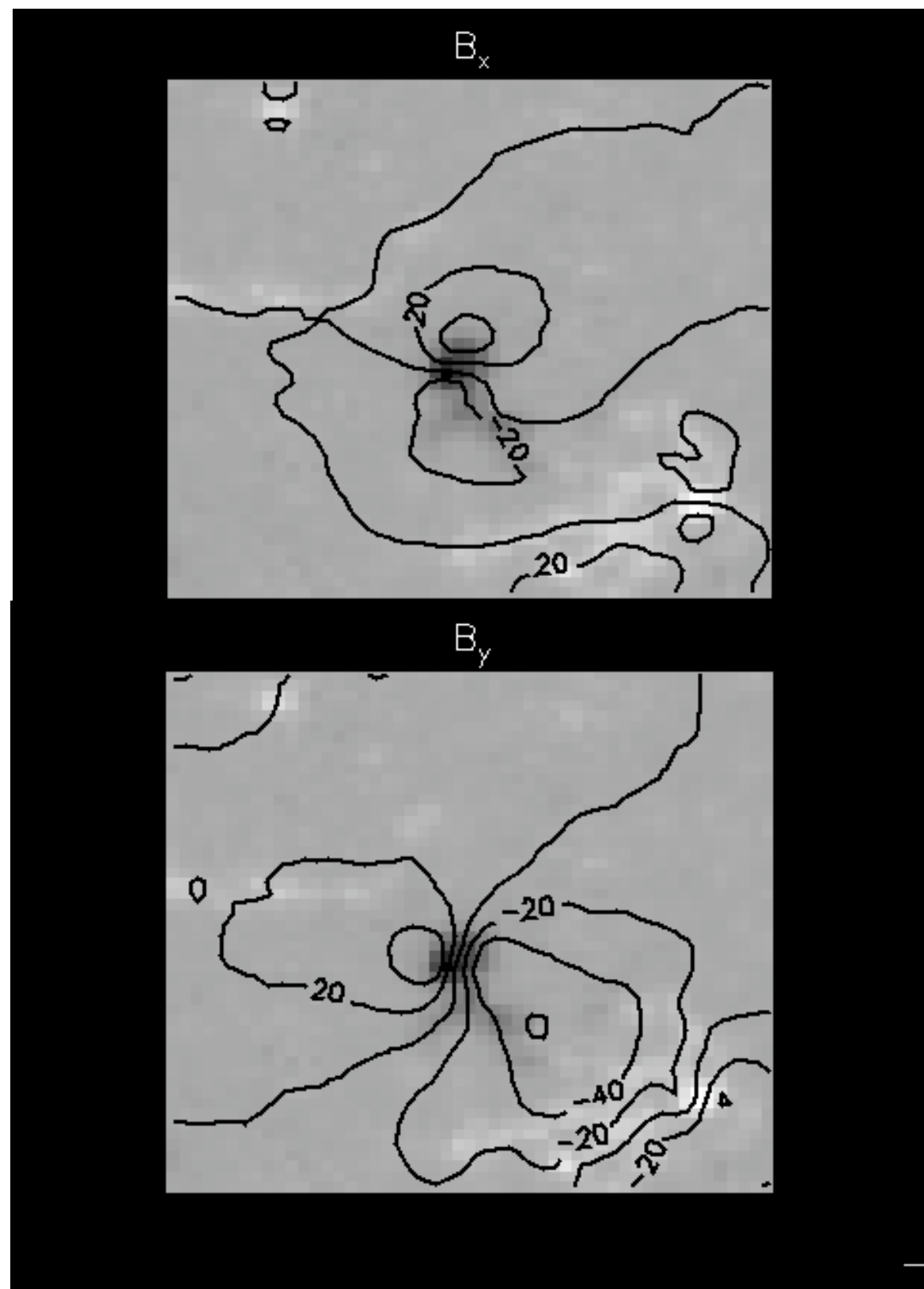
# Small AR, pores: including the chromosphere



detailed study of IBIS data: G. Cauzzi et al 2008, A+A

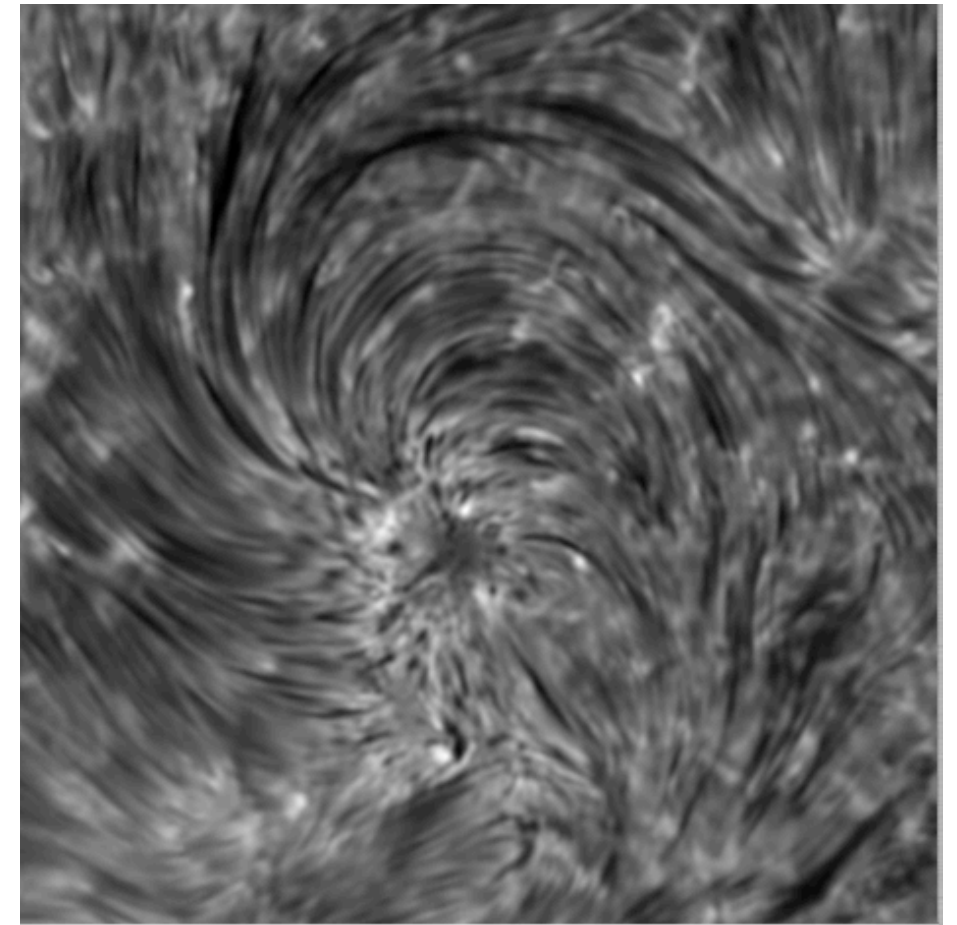
# Differences between potential and constant $\alpha$ photospheric fields

- IBIS morphology  $\Rightarrow$  transverse fields differ by  $\sim 20$ - $60$ G
- **Hinode** photospheric 630.2 sensitivity  $B_T(\text{app})$  Lites et al (2008) ApJ **672**, 1237
  - $40 \text{ Mx cm}^{-2} \text{ px}^{-1}$  (normal map)
- **current instruments can be used to study the forced  $\rightarrow$  force-free transition**
  - chromospheric electrodynamics

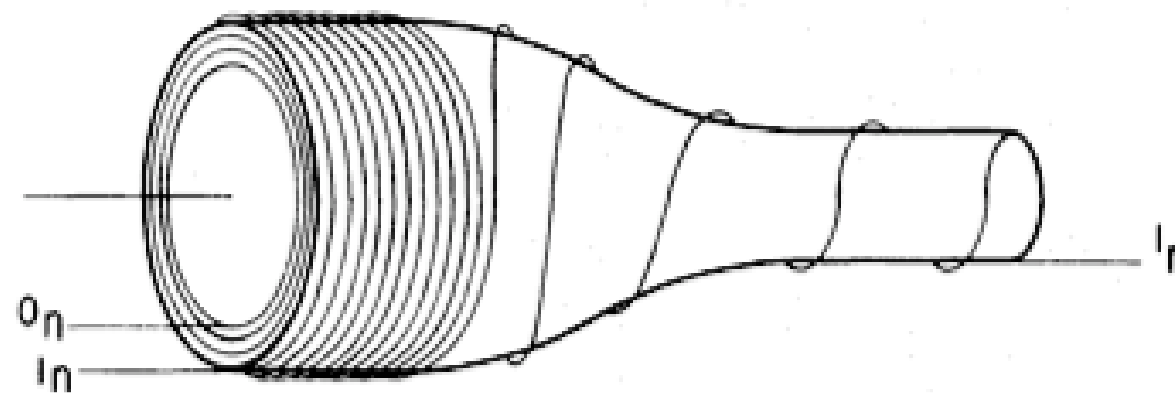


# Note: twist/ electrical currents can be easier to detect in the chromosphere!

- IBIS fibrils reveal a **clear  $B_\phi \Rightarrow j_z$**
- also Hinode rotating spicules
- Parker (1974):  **$B_\phi/B_z$  increases with  $z$**



## DYNAMICAL PROPERTIES OF MAGNETIC FIELD





**magnetic interface**  
**physical considerations**

# chromosphere as a partially ionized plasma

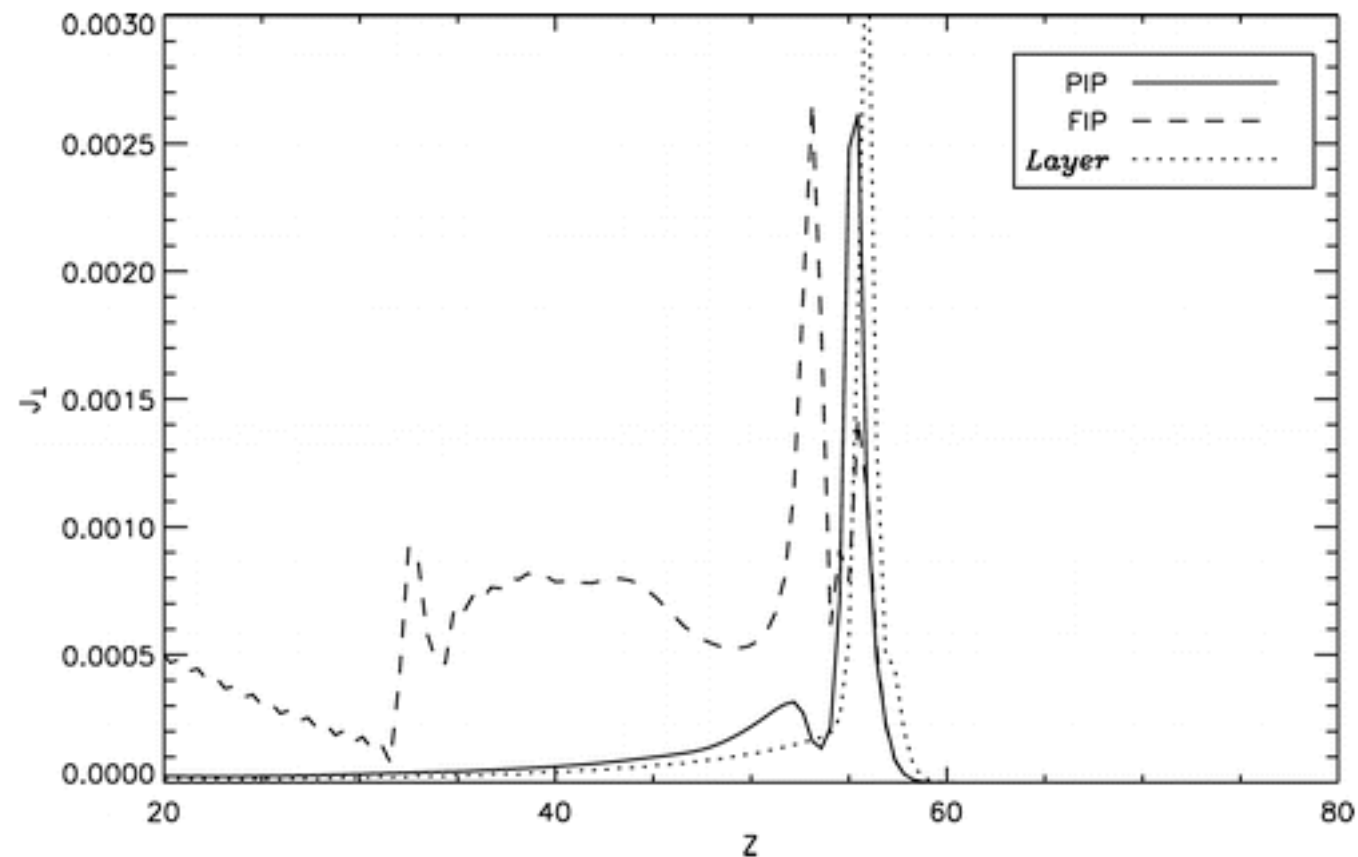
- partial ioniz<sup>n</sup> ⇒ 3-fluid *frictional dissipation, heating*
- efficient damping by ion-neutral collisions
- Kinetic theory ([Braginskii 1965](#))
  - $Q_{\text{fr}} = \mathbf{j} \cdot \mathbf{E} = j^2 / \sigma + (\xi_n \mathbf{j} \times \mathbf{B} - \mathbf{G})^2 / \alpha_n, \quad \mathbf{G} = \xi_n \nabla p - \nabla p_n$
  - “ambipolar diffusion”/star formation (1950s [Schlüter, Cowling](#))
- $\mathbf{G} = \mathbf{0} \Rightarrow$  “Cowling conductivity”  $\sigma_{\perp}^*$ 
  - $Q_{\text{fr}} = j_{\parallel}^2 / \sigma + j_{\perp}^2 / \sigma_{\perp}^* \quad \sigma / \sigma_{\perp}^* = 1 + 2 \xi_n \omega_e \tau_e \omega_i \tau_i, \quad \gg 1$
  - $\Rightarrow$  *rapid dissipation of  $\mathbf{j}_{\perp}$*
  - [Goodman & colleagues](#): wave heating
  - [Arber & colleagues](#): flux emergence

# Chromospheric dissipation of $\mathbf{j}_\perp$

- Braginskii (1965): certain motions ( $\mathbf{G}...$ ) dissipate  $\mathbf{j}_\perp$ 
  - Alfvén, fast modes, dynamic situations where
$$\nabla p - \rho \mathbf{g} + \mathbf{j} \times \mathbf{B} \neq \mathbf{0}$$
- **Not** slow modes, slow dynamics (cf. Goodman 2000)
- So, at coronal lower boundary, chromosphere makes:
  - $\mathbf{j}_\perp \sim \mathbf{0}$ ;  $\mathbf{j} \times \mathbf{B} \sim \mathbf{0}$
  - **weaker Alfvén/fast modes**

*Flux emergence: Arber, Haynes & Leake (2007) based upon Cowling's conductivity ( $\mathbf{G}=\mathbf{0}$ ):*

Plot of the magnitude of  $j_\perp$  as a function of height along the line  $x = y = 0$  for all three resistivity models at  $t = 160$ .



*...radical effect on  $\mathbf{j}$  and flux emergence process*

# chromosphere as a partially ionized plasma II

- $\sigma_{\perp}^*$  is some steps removed from  $\sigma$  (kinetic theory)
  - case  $\mathbf{G} \neq \mathbf{0}$ :  $\sigma_{\perp}^*$  incorrect!
  - one must consistently determine the nature of  $\mathbf{j}_{\perp}$  (cf. E-region electrojet) from the dynamics
- Fontenla (2005, 2008 A+A)
  - for length scales  $> 100$  km (few mHz waves),
  - $Q_{\text{fr}} = \mathbf{j} \cdot \mathbf{E}$  too small, invokes instability (Farley-Buneman)
  - need neutral component velocity  $>$  ion acoustic velocity



# **thermal interface**

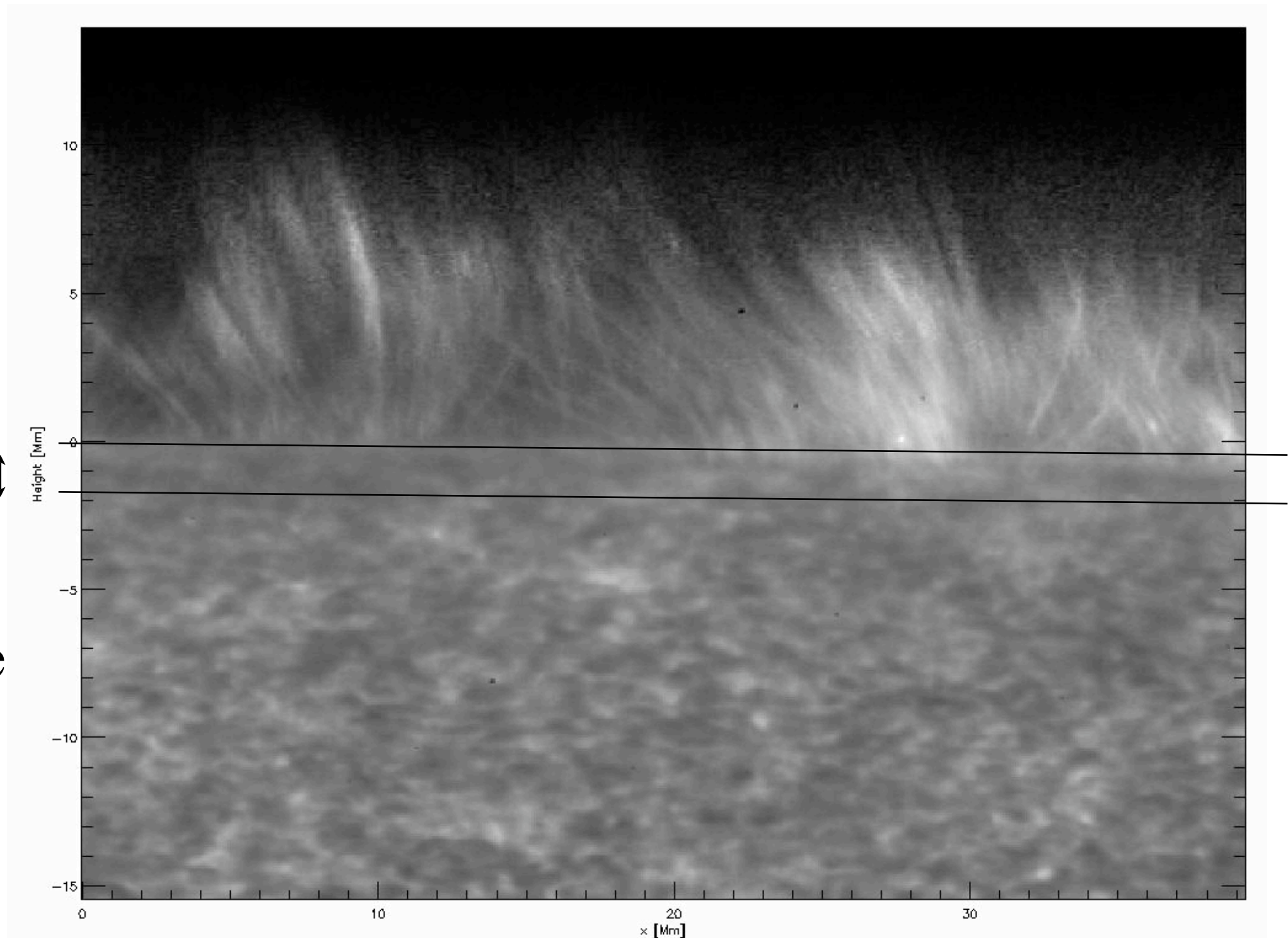


# non-planar interface: Hinode spicules

- Ca II (radial filter to enhance spicules, M. Carlsson)

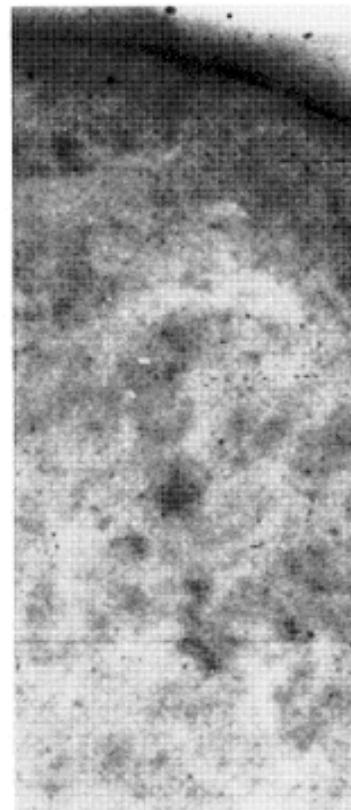
nb.  
stratified VAL  
chromosphere  
1.5Mm only

spicules are  
produced by the  
chromosphere



# non-planar interface: transition region-corona

- Feldman and colleagues
  - different morphology
  - TR thermally, magnetically
  - radiating entity = “



Mg I



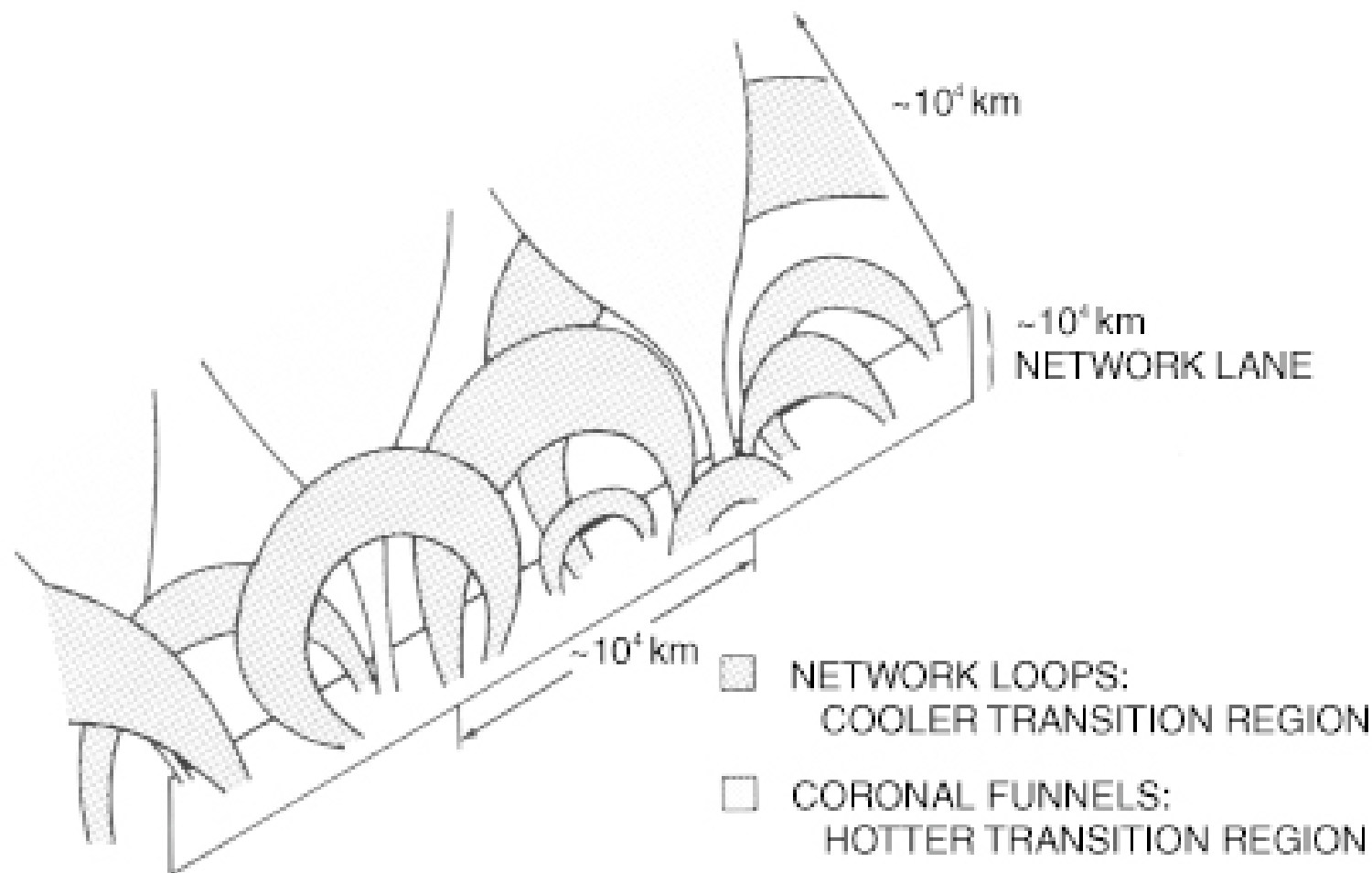
*what is the nature of  
this emitting plasma?*

C III (977 Å)



## Dowdy et al. (1986)

- Mixed polarity within network boundaries
- tries to explain “UFS”
- indeed these are thermally and magnetically separate entities



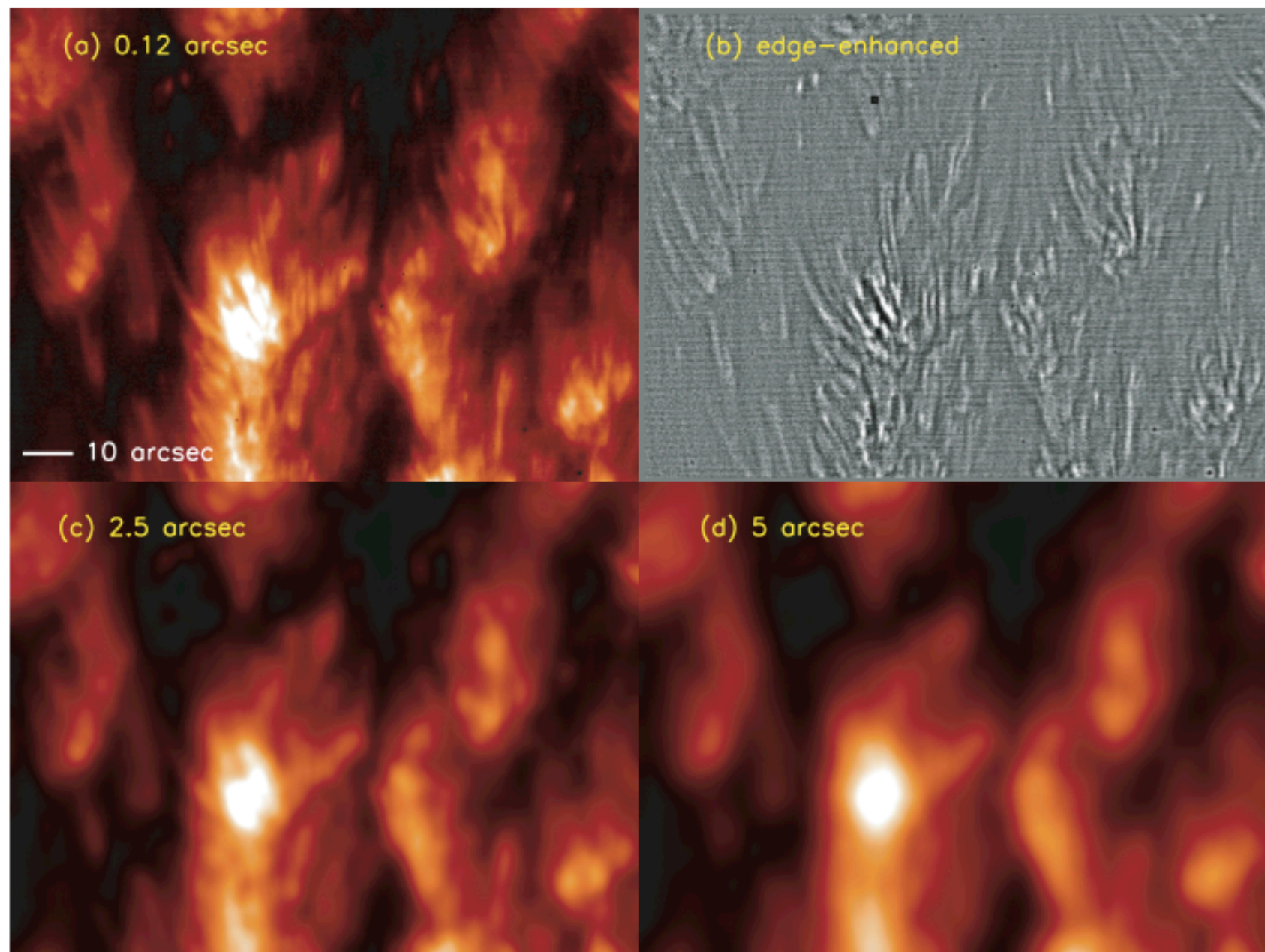
Questions: stability; footpoint magnetic fields; fate of downward conductive energy flux,...



# Judge & Centeno (2008)

- VAULT  $L\alpha$  data vs. KPNO magnetic data
  - supplemented by Hinode SP vector polarimetry
- Prompted by Patsourakos et al (2007)
  - We noted something “odd” about proposed cool loops
  - **large-scale alignment of  $L\alpha$  threads**

Patsourakos et al:





# KPVT+POTL FIELDS+VAULT active network

Potential fields:

Black=low loops ( $h < 5\text{Mm}$ )

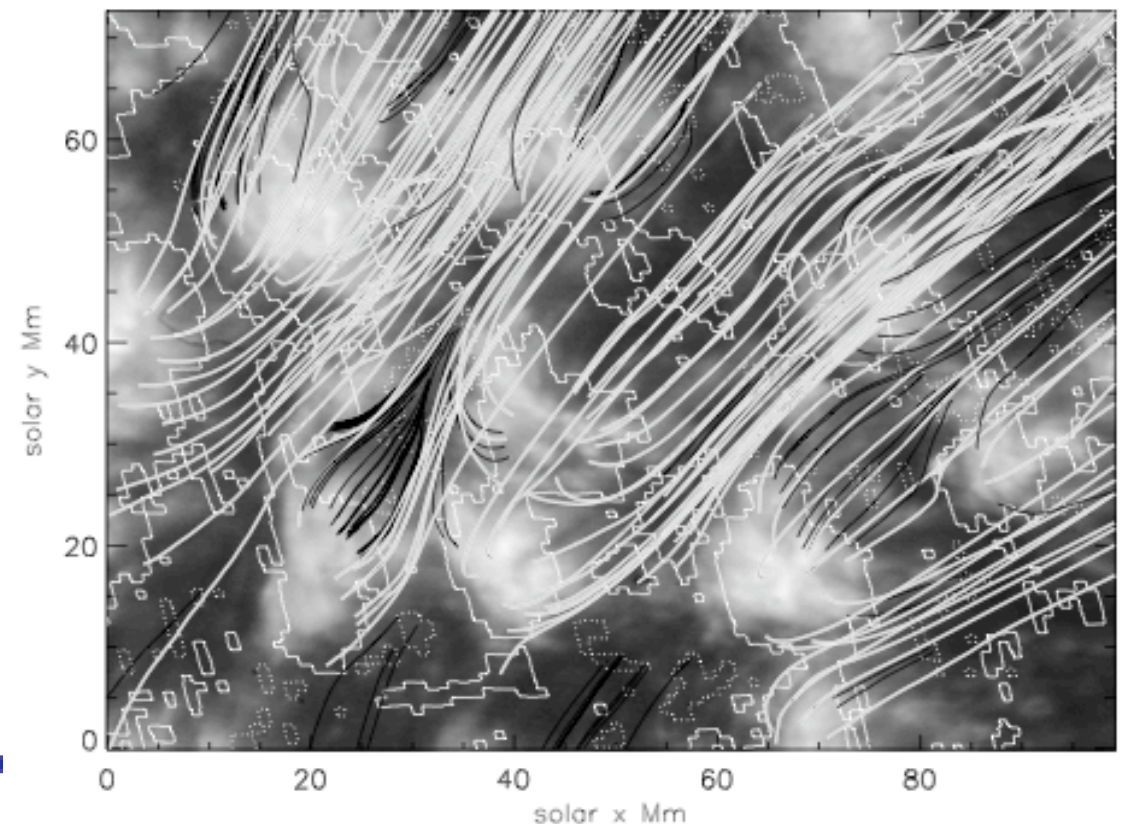
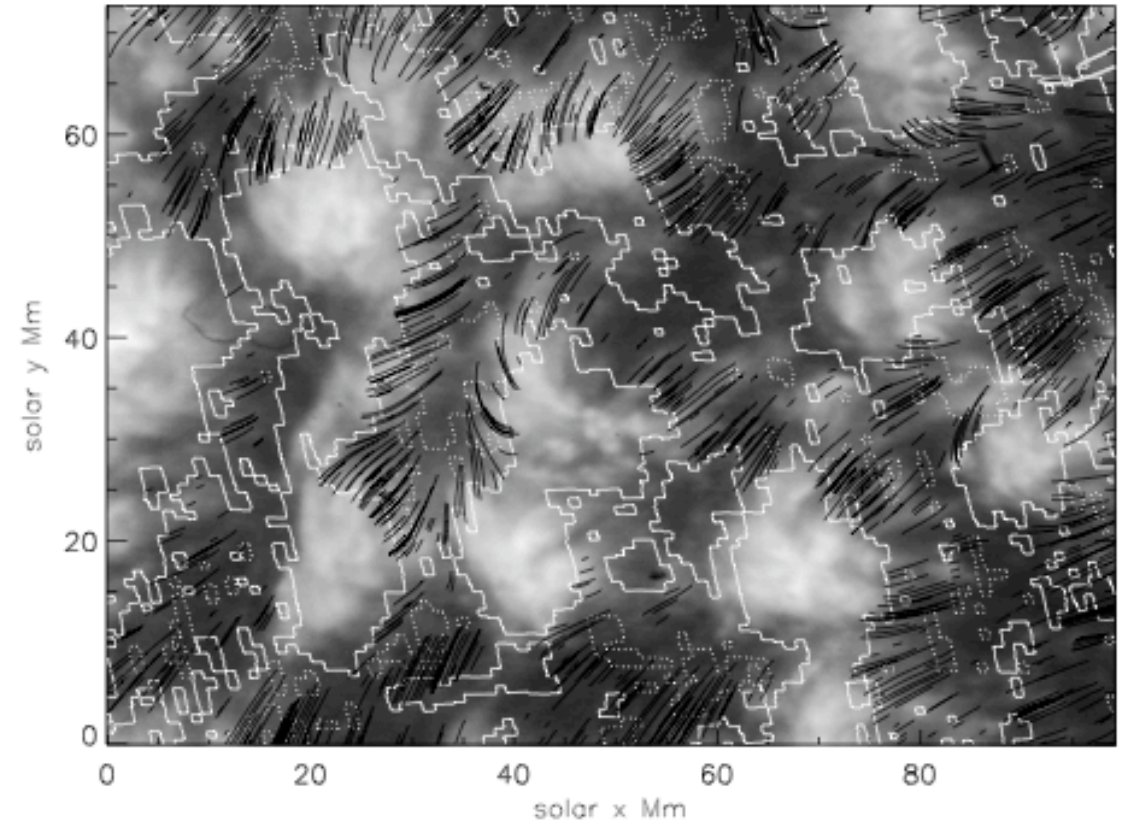
Gray= higher

Non potential fields are generally higher

Stability requires that higher loops be hot

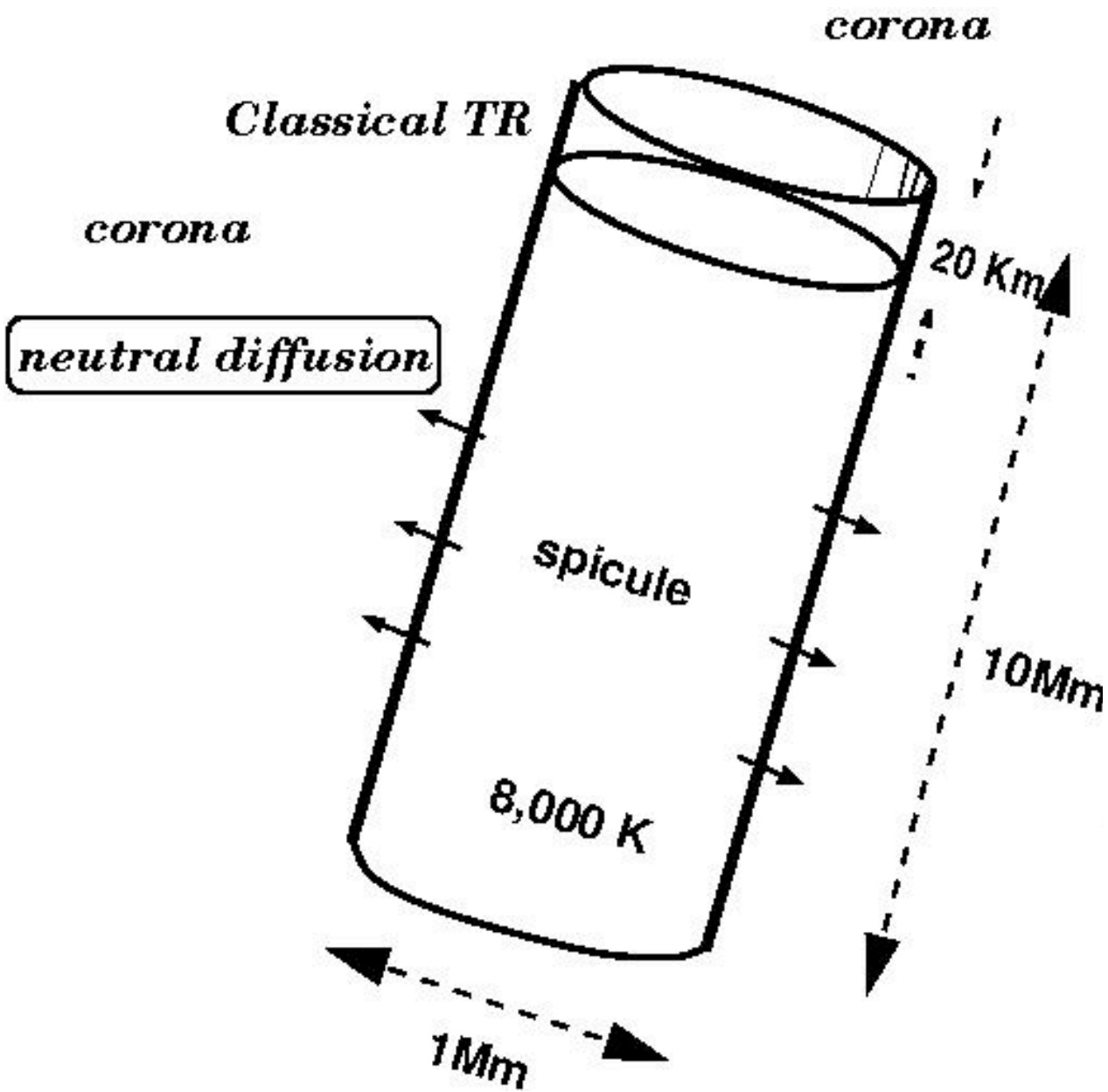
Most  $L\alpha$  emission originates from the base of hot, coronal loops

Unresolved mixed polarity fields don't work.  $L\alpha$  emission forms above  $h=0.8$  Mm. Loops with footpoints separated by  $\leq 1''$  cannot reach these heights



**Judge (2008) ApJL 683, 87-90:  
is there a simple explanation for the network  
transition region without appealing to cool  
loops?**

# “spicule” → cross field diffusion → TR radiation



Initial corona

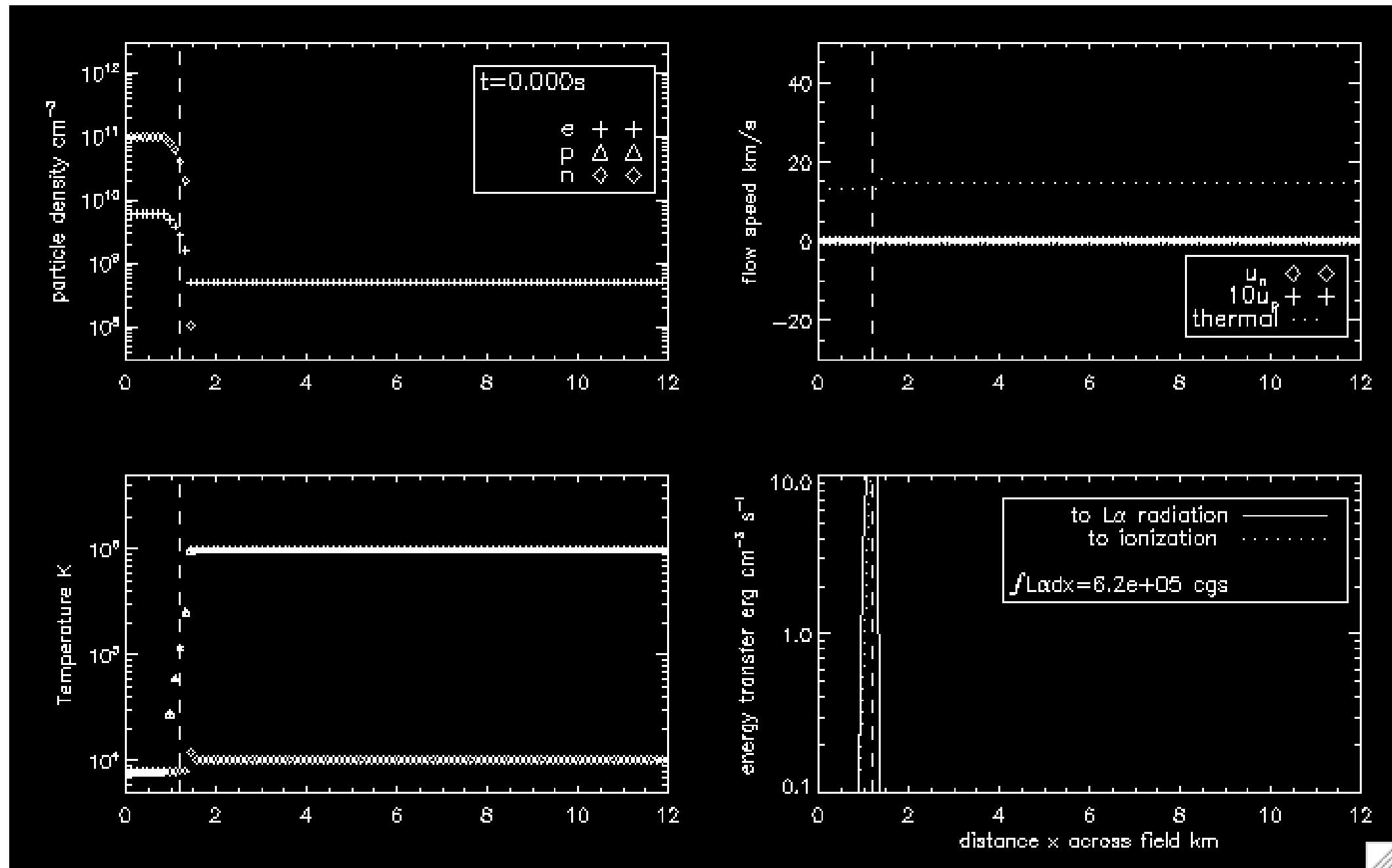
$T_h$	K	$10^6$	
$n_h$	$\text{cm}^{-3}$	$8.0 \times 10^8$	
$n_p, n_e$	$\text{cm}^{-3}$	$4.0 \times 10^8$	
$p$	$\text{cm}^{-3}$	$1.1 \times 10^{-1}$	
$B$	G	10	
$\beta$		$2.8 \times 10^{-2}$	
$\omega_p$	$\text{s}^{-1}$	$9.6 \times 10^4$	
$\tau_{gyro}$	km	$1.5 \times 10^{-3}$	
$\tau_{pp}$	s	1.6	$n_p^{-1} T^{+3/2}$
$\omega_p \tau_{pp}$		$1.5 \times 10^5$	
$\tau_{ee}$	s	$5.0 \times 10^{-2}$	$n_e^{-1} T^{3/2}$
chromospheric tube			
$T_c$	K	$8.0 \times 10^3$	
$\bar{v}$	$\text{km s}^{-1}$	13	$T^{1/2}$
$n_c$	$\text{cm}^{-3}$	$10^{11}$	
$\tau_{nn}$	s	$1.4 \times 10^{-2}$	$n_n^{-1} T^{-1/2}$



# Results: model $L\alpha \sim 0.1x$ observed using only local coronal heat

1D 3-fluid  
calculation  
of cross-field  
diffusion  
from a cool  
flux tube into  
coronal  
plasma

no field  
aligned  
conduction



calculations with different coronal  $n, T$ : non-linear  
relationship between  $L\alpha$  and coronal emission

# Judge (2008)

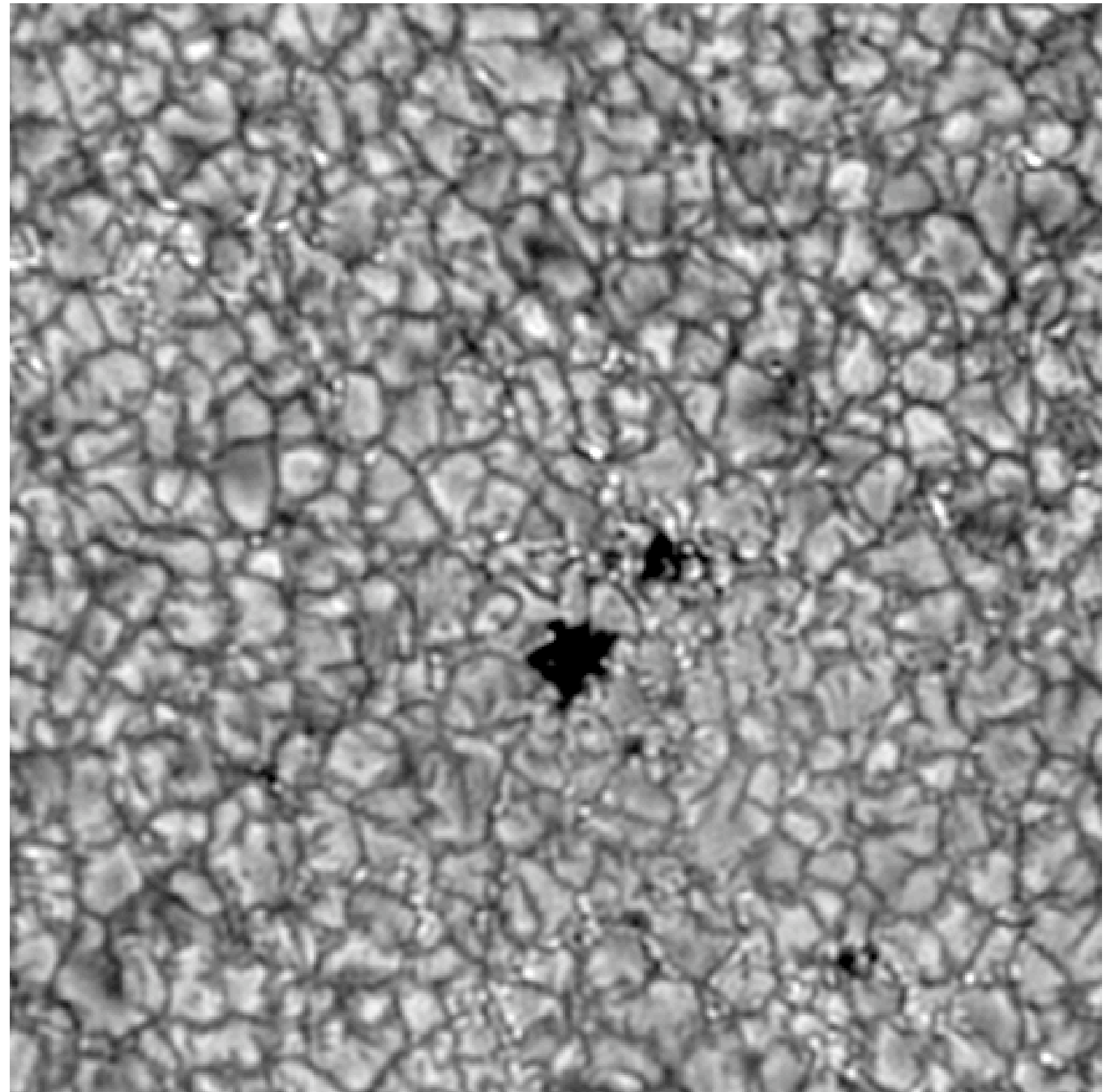
- chromosphere supplies the (neutral) mass, corona the energy
  - “UFS” in this new picture is **thermally connected to the corona**
  - calculations for  $L\alpha$  are promising, (also  $L\beta$ , He I 584)
- **cross-field diffusion of neutrals** might solve the 40+ yr problem of **energy balance in extended structures in the lower TR**
- needed
  - 2D calculations including field-aligned conduction and dynamics
  - **observations of the chromosphere/corona interface in relation to magnetic field**

# **chromosphere as the coronal base**



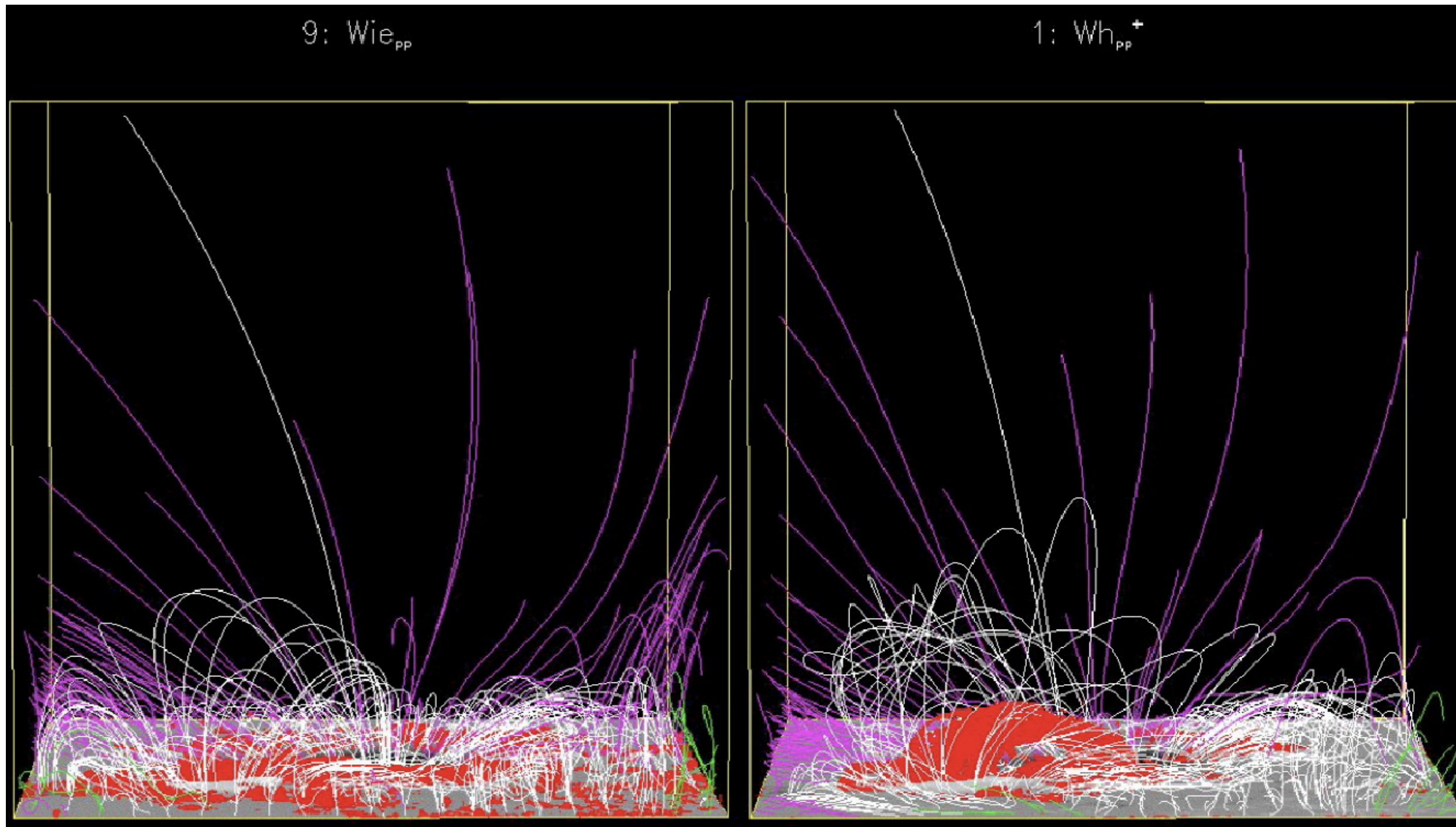
# the need for sub-arcsec imaging spectroscopy/ spectropolarimetry at the coronal base

- to map low  $\beta$  magnetic structure
  - $\leq 1''$  angular resolution
  - $\geq 30'' \times 30''$  FOV
- to get dynamics
  - $< 30$ s (spicules..)
- ground-based:
  - spectropolarimetry
  - Ca II, He I, H I,...
  - chromosphere only
- space (UV)
  - spectroscopy
  - chrom+TR+corona



# nlff field extrapolation (Schrijver et al 2008)

red:  
current



new chromospheric **B** constraints (including e.g., just fibrils) can provide boundary conditions compatible with the calculations



# To understand the corona we must study what is under Gold's line

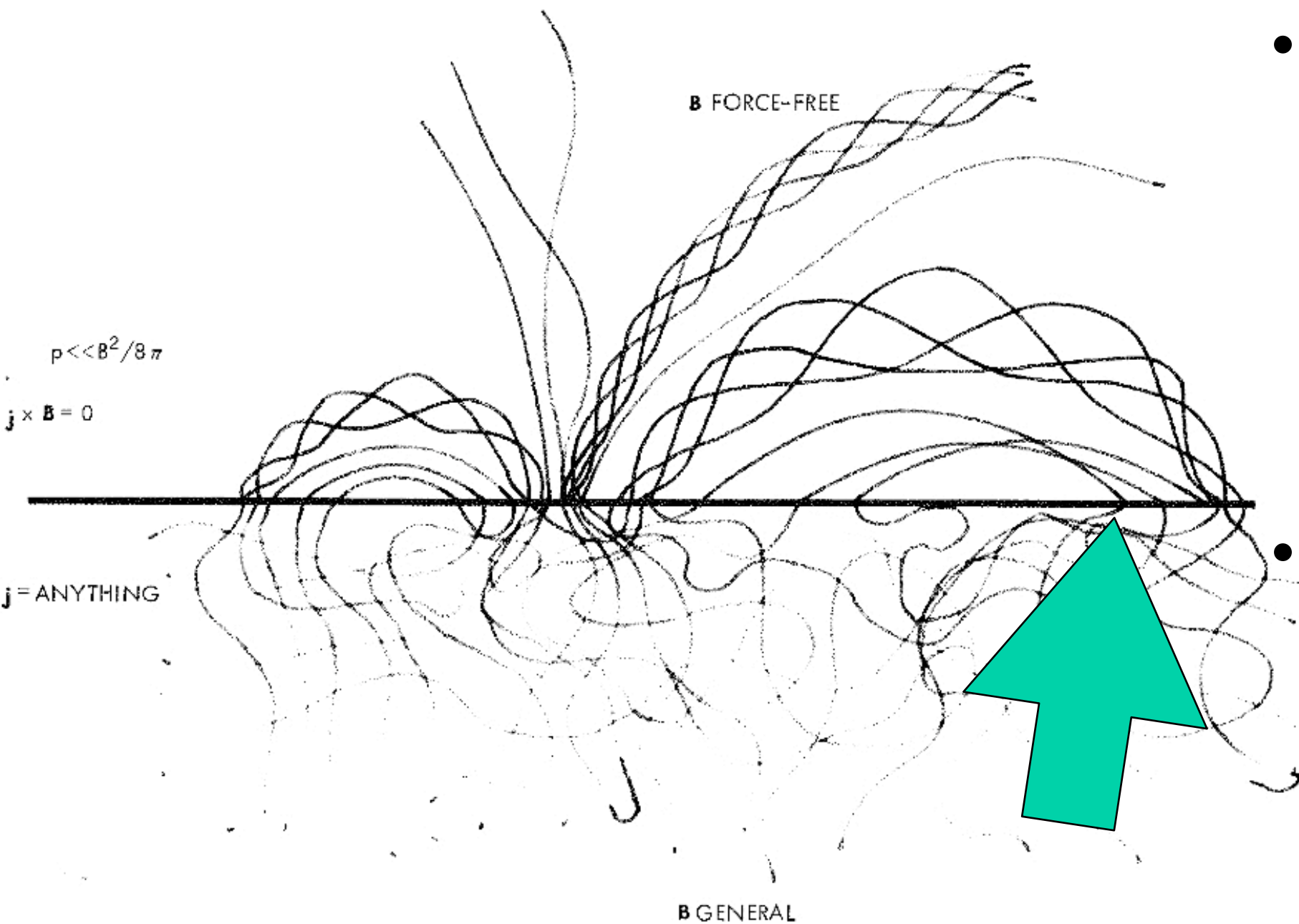


FIGURE 44-2. Magnetic field in a turbulent conducting medium. The fluid pressure is assumed large compared with magnetic forces below the dividing plane and small above it.

- understand mass, momentum, free energy transport across chromosphere
  - *single fluid MHD OK?*
- obtain much better boundary conditions on the corona
  - magnetic
  - thermodynamic
- *ground+space data required*