

# The photosphere/corona interface: new perspectives

**Philip Judge** + R. Centeno, M. Kubo, B. Lites, S. McIntosh, A. G. de Wijn, **HAO**; G. Cauzzi, K. Reardon, **Arcetri**; A. Tritschler, H. Uitenbroek **NSO**

*re-visiting some physical issues*

old vs. new perspectives

magnetic interface

thermal interface

September 2008

PLATE X

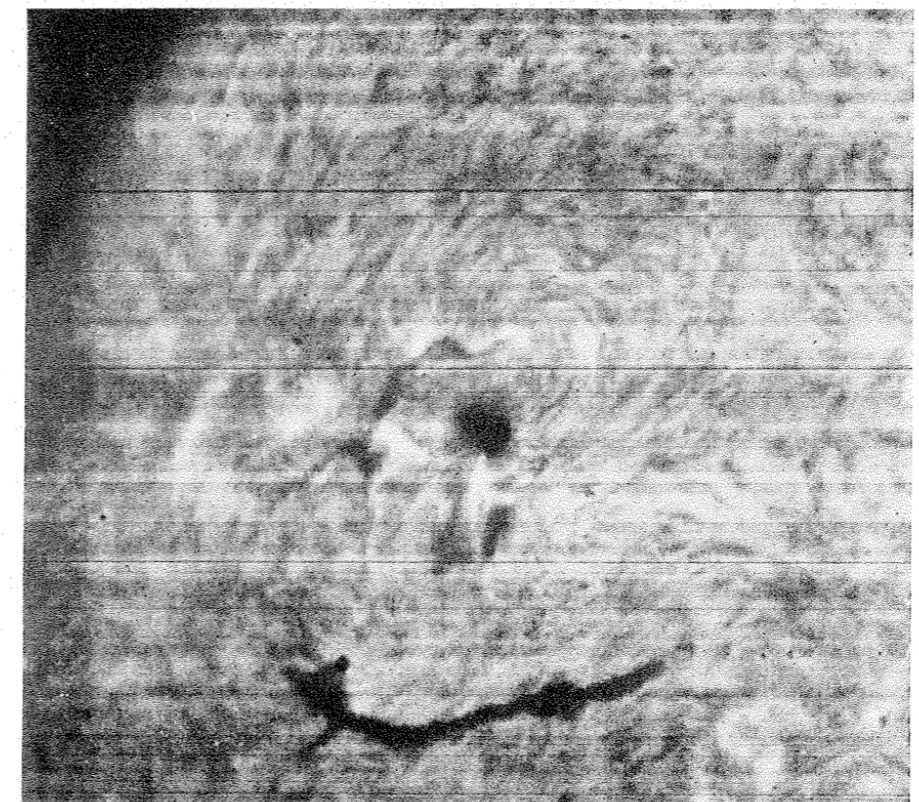


FIG. 1.—SUN-SPOT AND HYDROGEN ( $H\alpha$ ) FLOCCULI  
1908, May 29, 4<sup>h</sup> 26<sup>m</sup> P. M. Scale: Sun's Diameter = 0.3 Meter



NCAR

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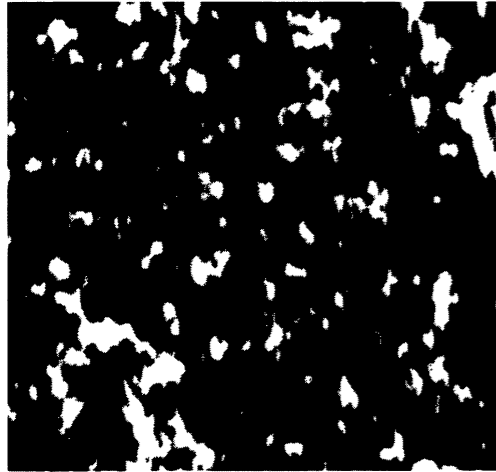


# the chromosphere

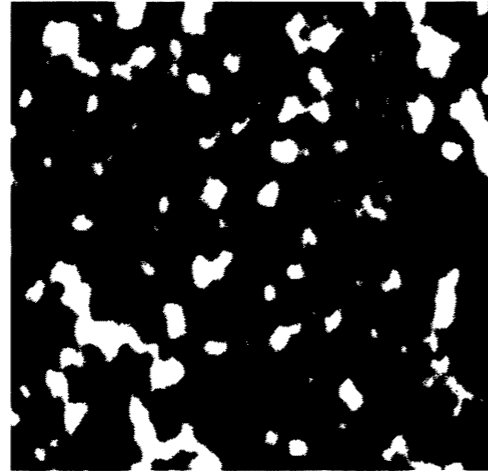
- stratified: spans 9 pressure scale heights
- requires 30-100x as much power as the corona
- usually contains plasma  $\beta=1$  surface
- is the lower boundary for the corona
  - modulates flow of mass, momentum, energy and magnetic field into the corona
  - implicit mass reservoir in coronal loop scaling laws
- yet
  - “chromosphere Hinode” search reveals 1/3 of “corona Hinode” publications
  - chromosphere is an “ignore-o-sphere”?
  - “too complicated”?

# Example of “old” perspectives SKYLAB data - VAL thermal models

1973 JULY 8

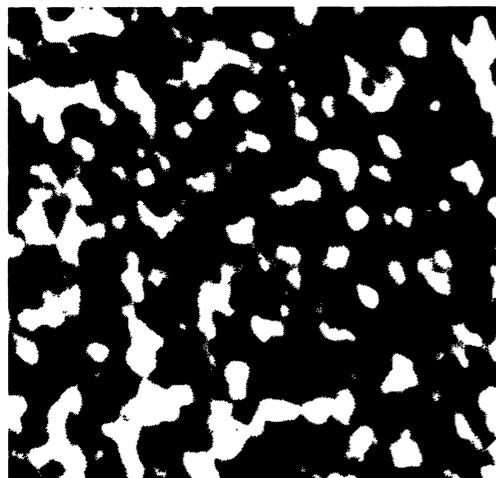


C I Continuum ( $\lambda=104.4\text{nm}$ )

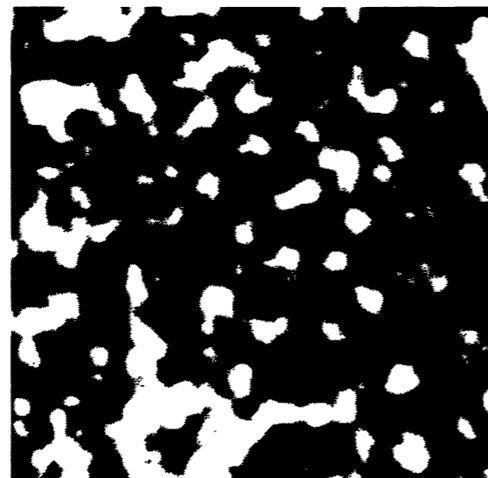


Lyman Continuum ( $\lambda=74.0\text{nm}$ )

1973 JULY 9



Lyman Alpha Wing ( $\lambda=121.1\text{nm}$ )



Lyman Continuum ( $\lambda=90.7\text{nm}$ )

1D  
RT  
nLTE  
HSE  
PRD

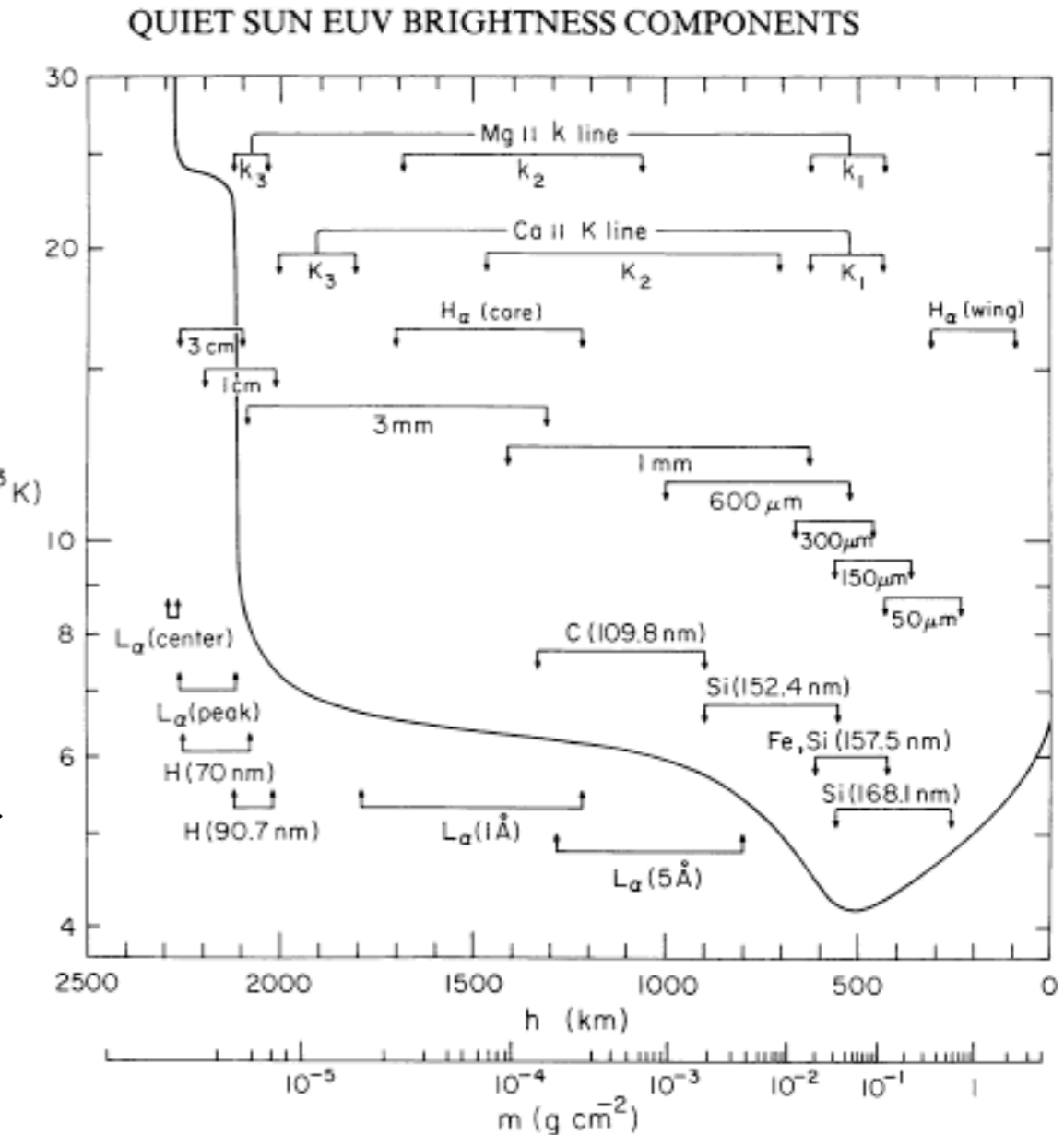
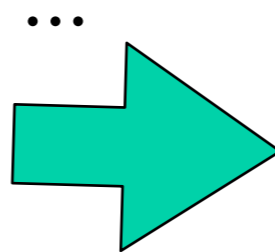


FIG. 6.—Skylab spectroheliograms of two  $5' \times 5'$  areas of the solar surface, at wavelengths 104.4 and 74 nm (above) and wavelengths 121.1 and 90.7 nm (below).

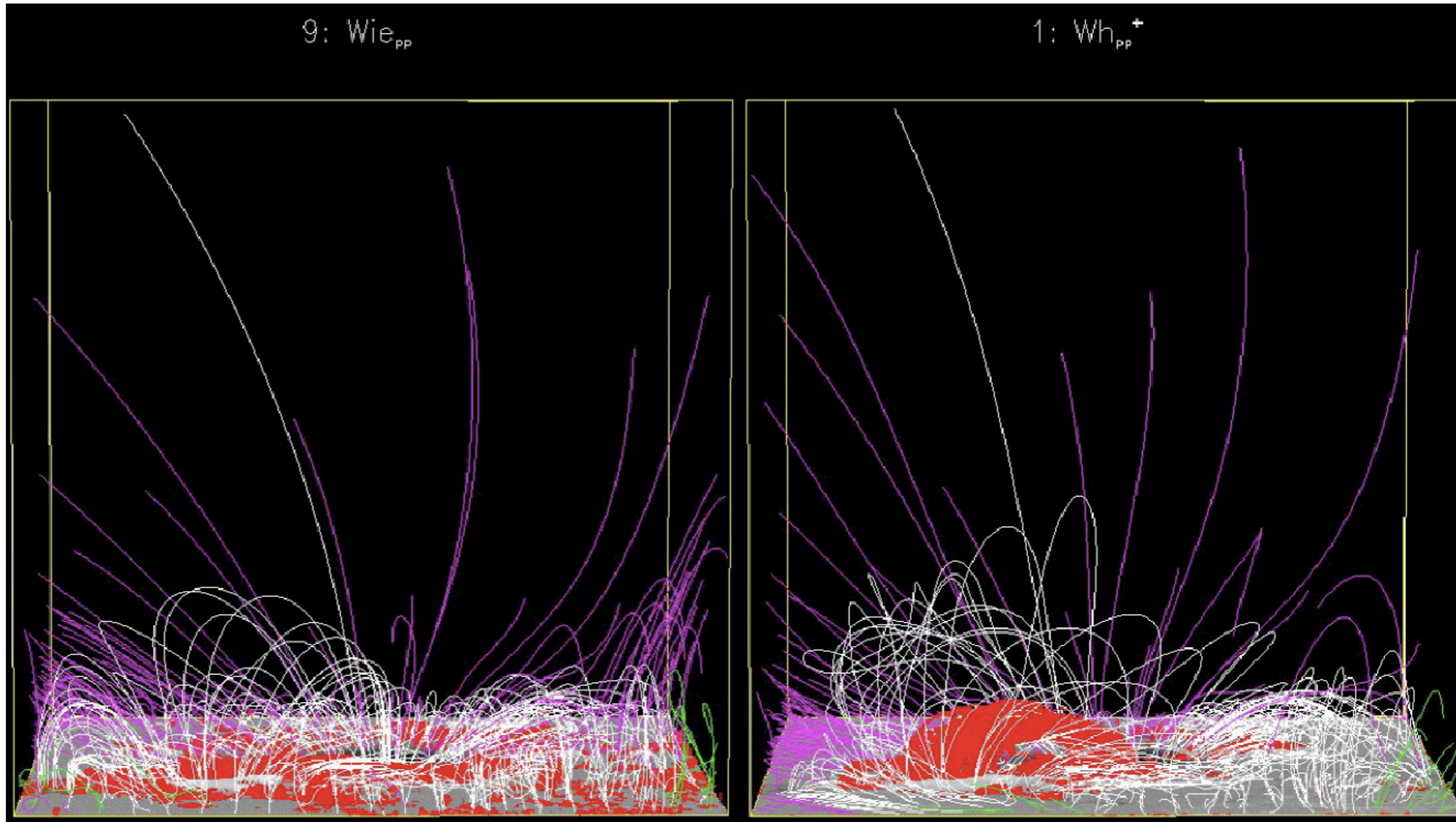
VERNAZZA *et al.* (see page 647)

Heroic reference work of vital importance, 1981



# Recent(!) example of “old” perspectives nlff field extrapolation (Schrijver et al 2008)

red:  
current



Hinode SP photospheric vector polarimetry, no  
chromospheric data (nb. Low & Flyer 2007)

# **New perspectives: DOT and TRACE**

**9 Jul 2005 (A.G. de Wijn, R. J. Rutten)**

photosphere

chromosphere

corona

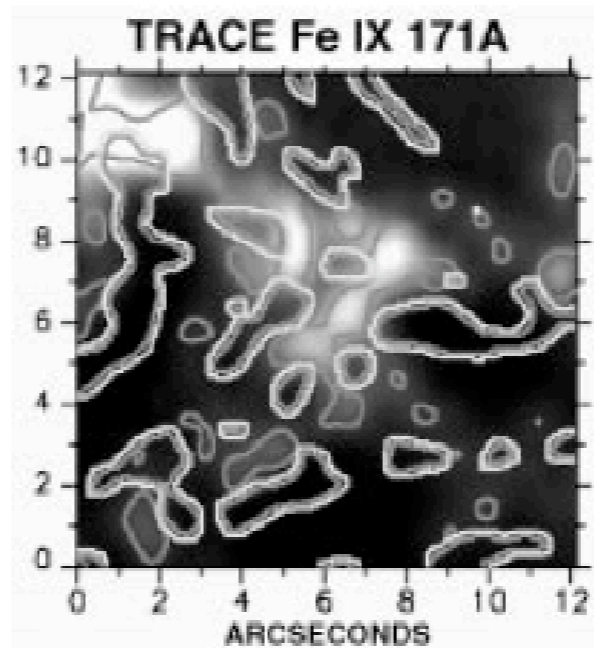


# **magnetic interface**

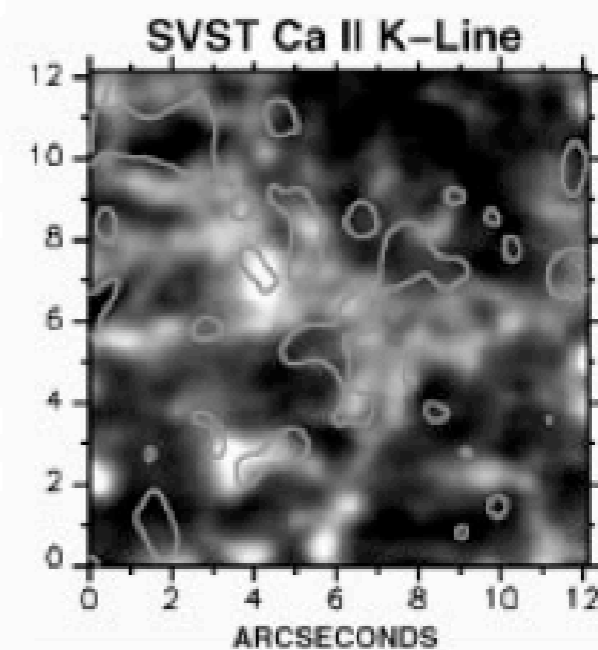
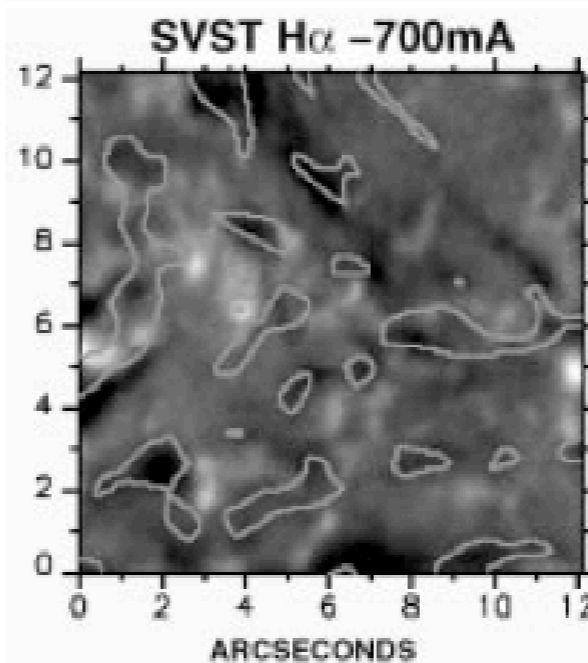
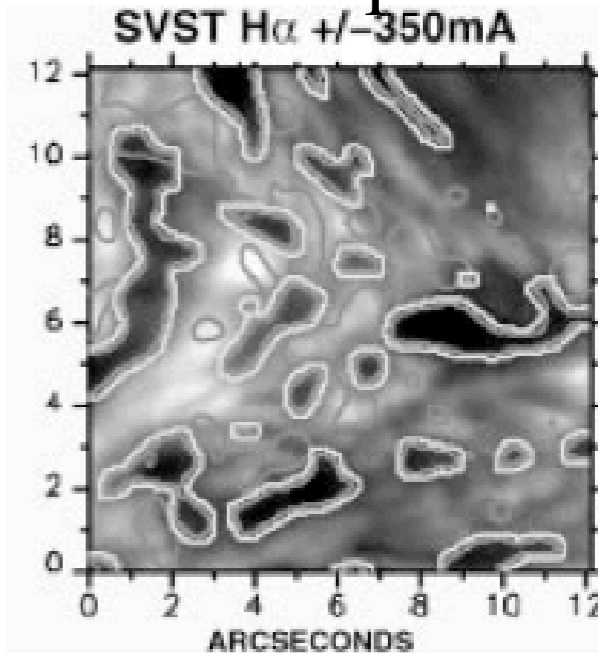
# Magnetism and the solar atmosphere

- measure  $B$  where possible
- high plasma conductivity-  
“trace field lines” from  
photosphere to corona
- TRACE & other missions  
failed to do this
- why?- chromosphere

lower corona



upper  
chromosphere



De Pontieu et al. 1999  
“moss”

upper  
chromosphere

magnetic elements +  
reverse granulation

## Gold (1964)

- consider potential and f-f fields in upper half
- the electro-dynamics of the **chromosphere** is **critical** to the supply of magnetic free energy into the corona.
- traditionally it is treated as in the figure

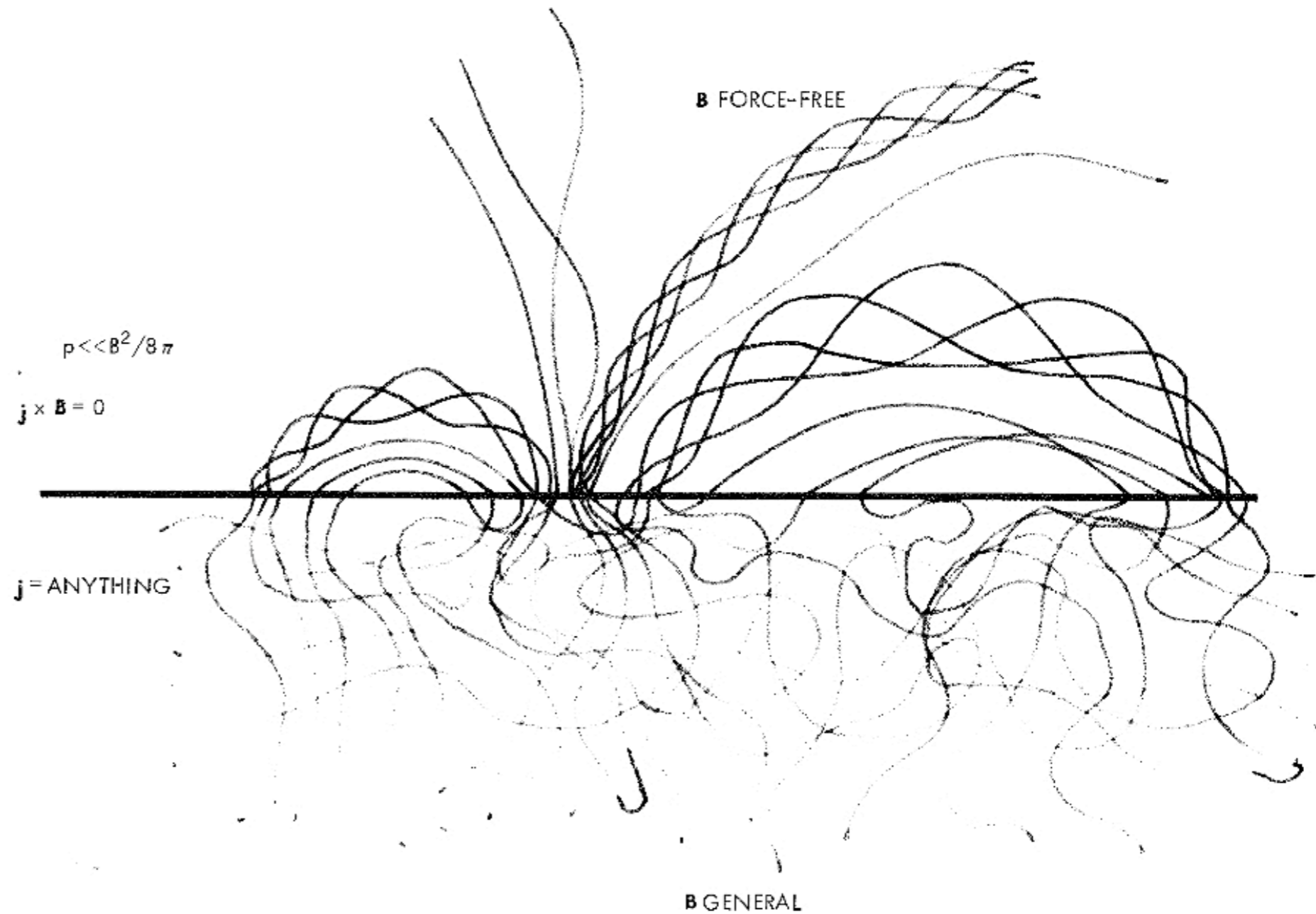
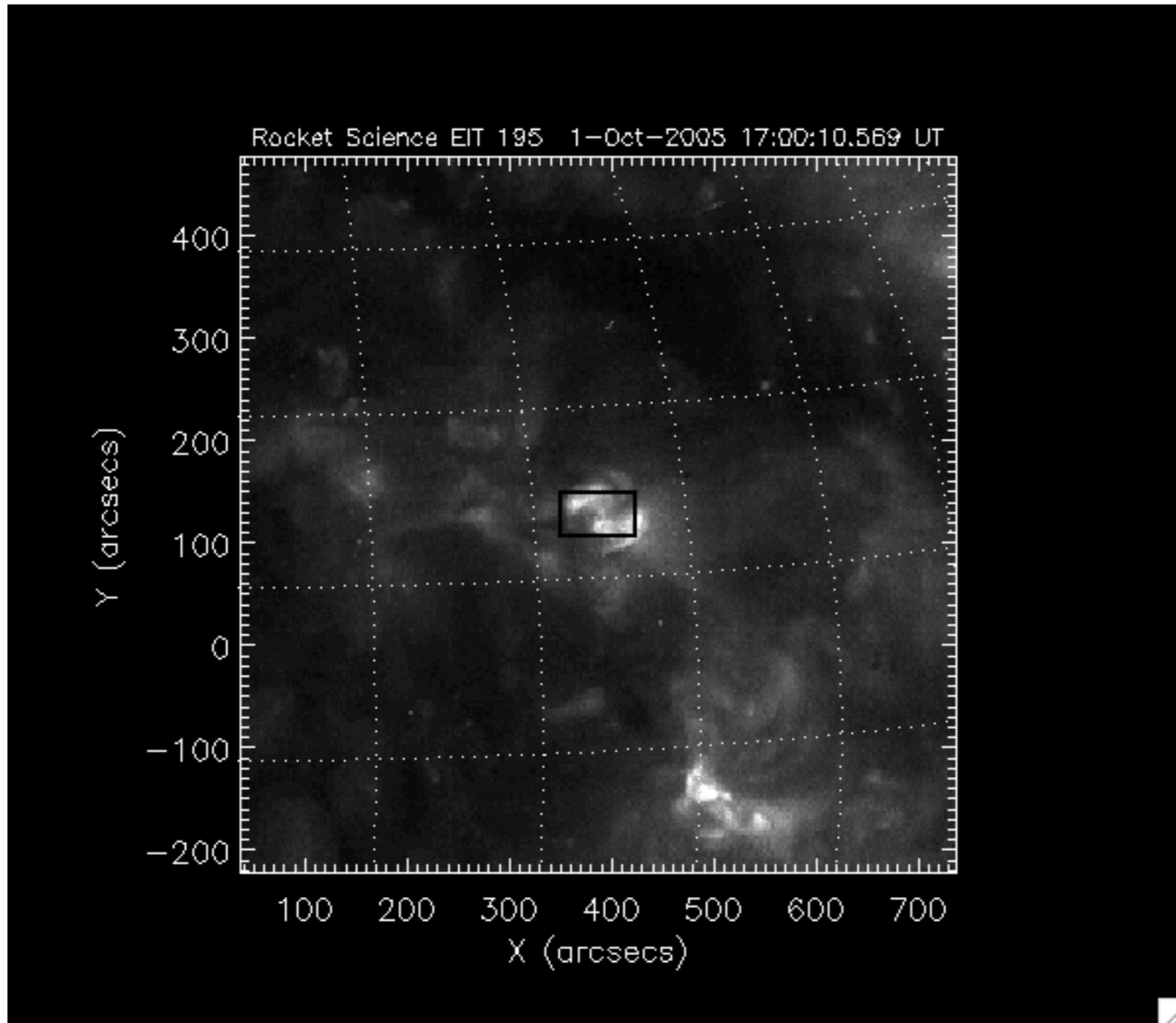


FIGURE 44-2. Magnetic field in a turbulent conducting medium. The fluid pressure is assumed large compared with magnetic forces below the dividing plane and small above it.

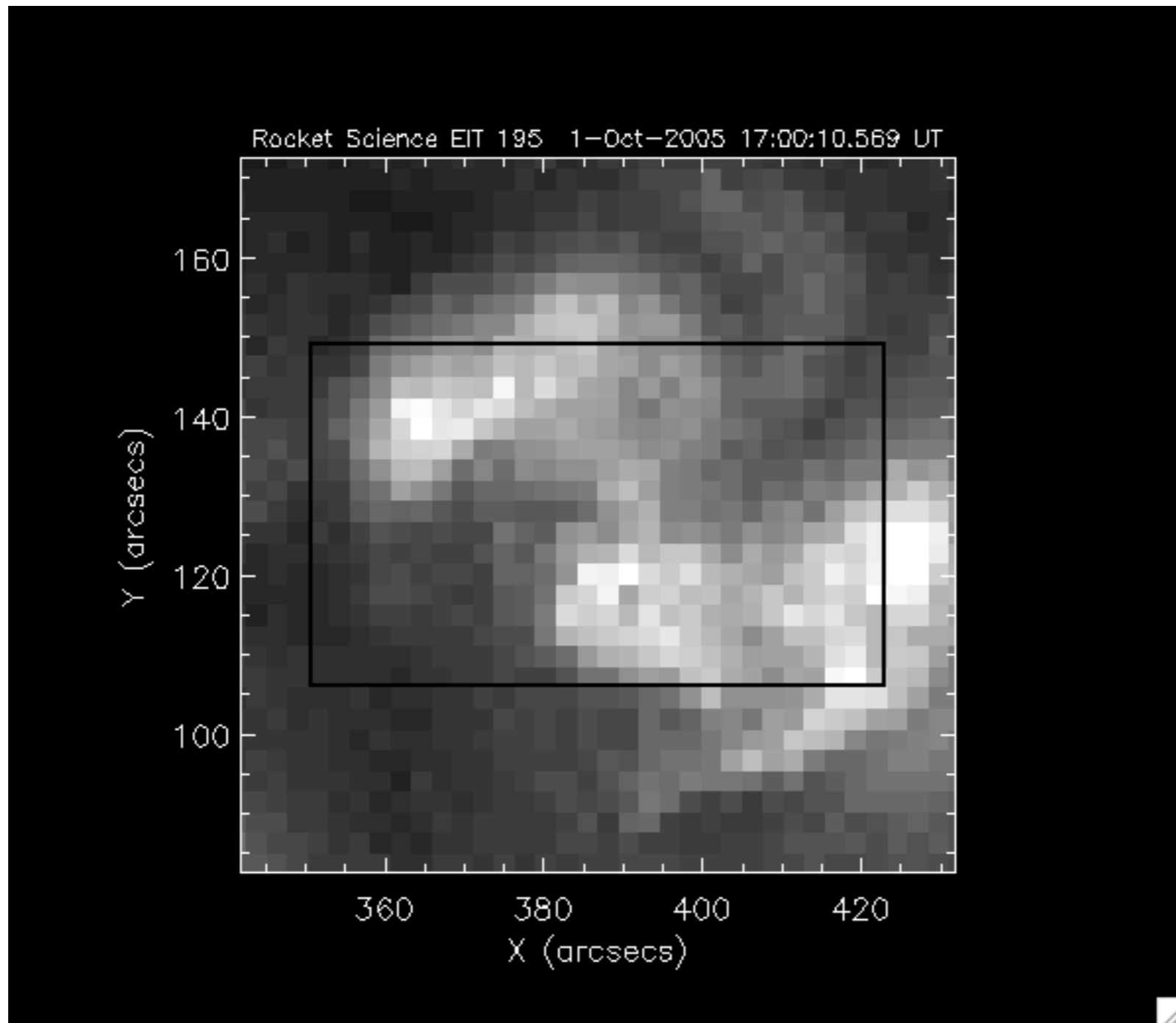


**magnetic interface  
observations:  
an example**

# Small AR, pores

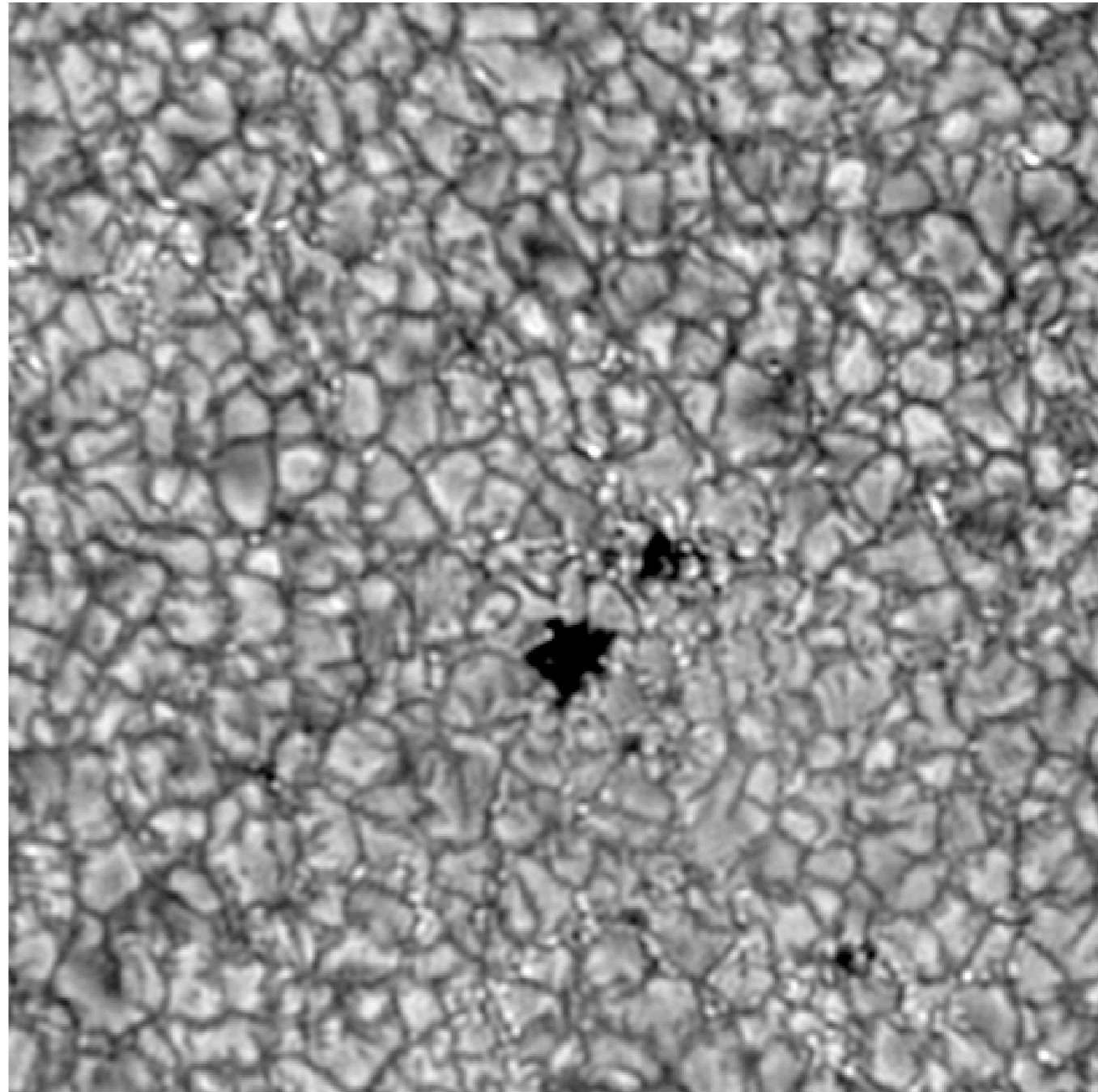


# Small AR, pores: closer view



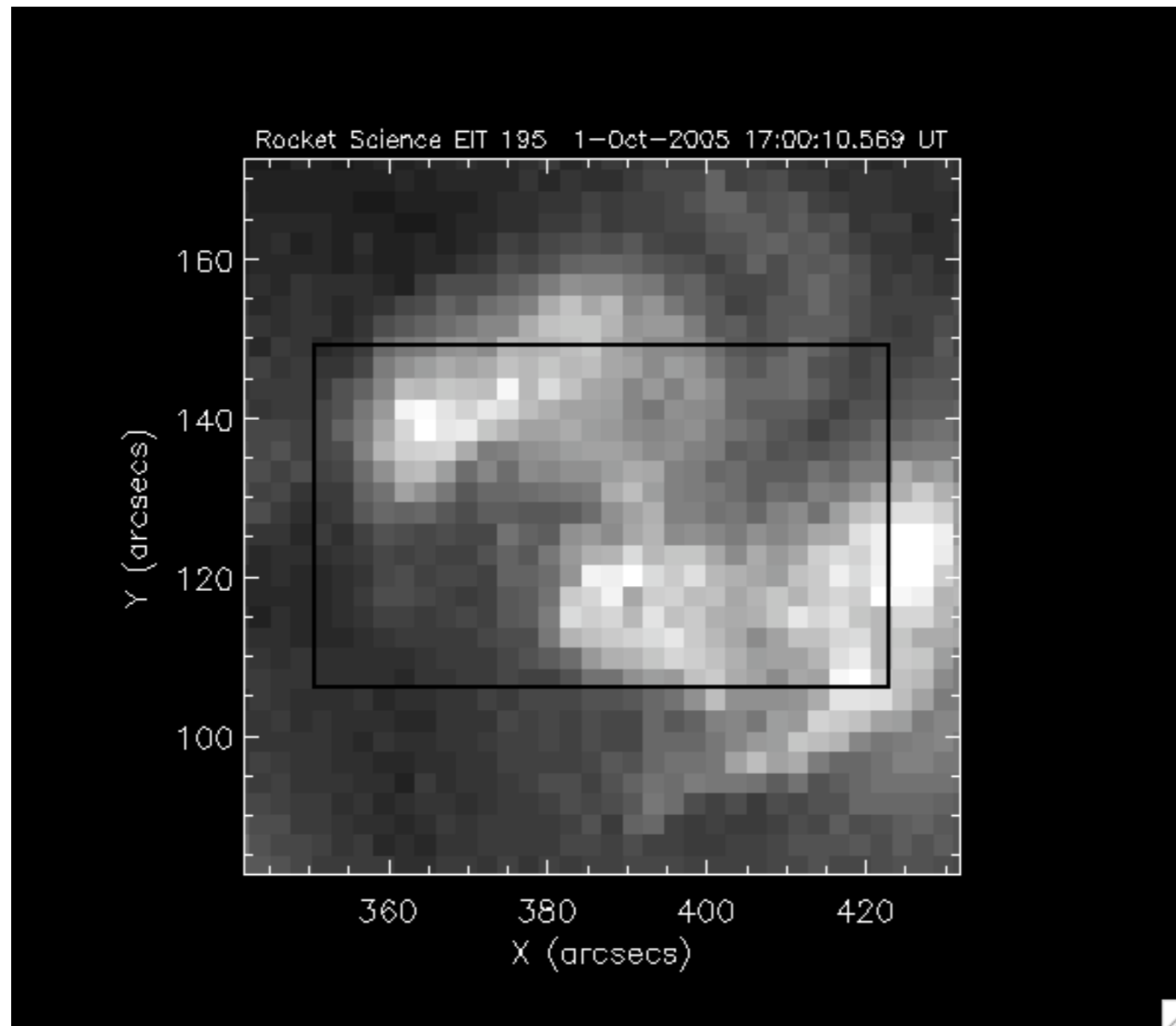
# Chromosphere as seen with IBIS

- Ca II 854.2 nm
- samples many pressure scale heights
- base of corona is **very** different from photosphere



G. Cauzzi et al 2008, A+A

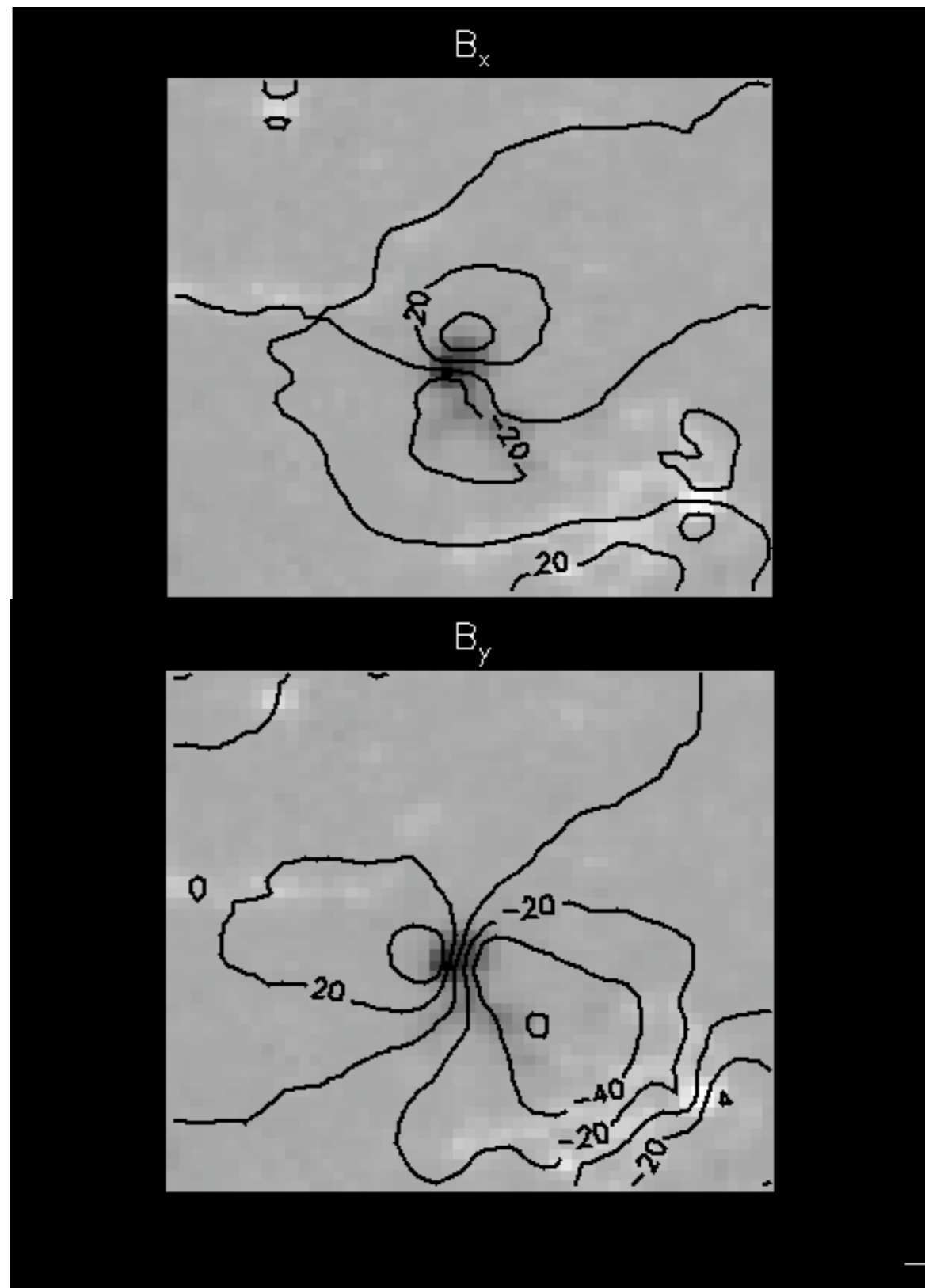
# Small AR, pores: high resolution photosphere and chromosphere



detailed study of IBIS data: G. Cauzzi et al 2008, A+A

# Differences between potential and constant $\alpha$ photospheric fields

- IBIS morphology  $\Rightarrow$  transverse fields differ by  $\sim 20\text{-}40\text{G}$
- **Hinode** 630.2 sensitivity  $B_T(\text{app})$   
Lites et al (2008) ApJ **672**, 1237
  - $40 \text{ Mx cm}^{-2} \text{ px}^{-1}$  (normal map)
  - $20 \text{ Mx cm}^{-2} \text{ px}^{-1}$  (deep map)
- **Hinode can study photospheric vs chromospheric electrical currents, forced  $\rightarrow$  force free transition!**
- Total  $\div$  potential energy:
  - 2 (chromosphere)
  - 5-10 (corona)



# Hale 1908: 100 years on

PLATE X

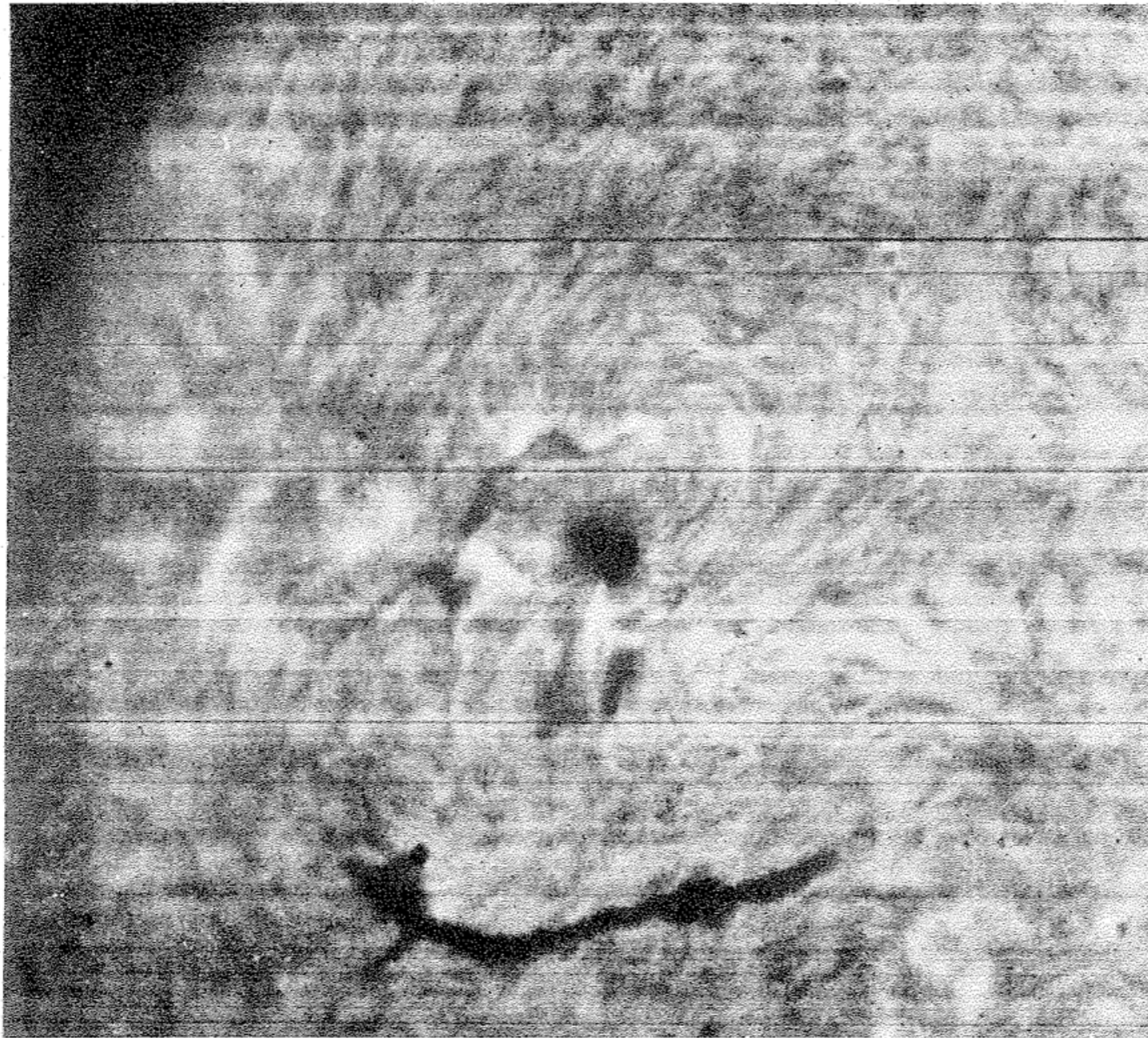


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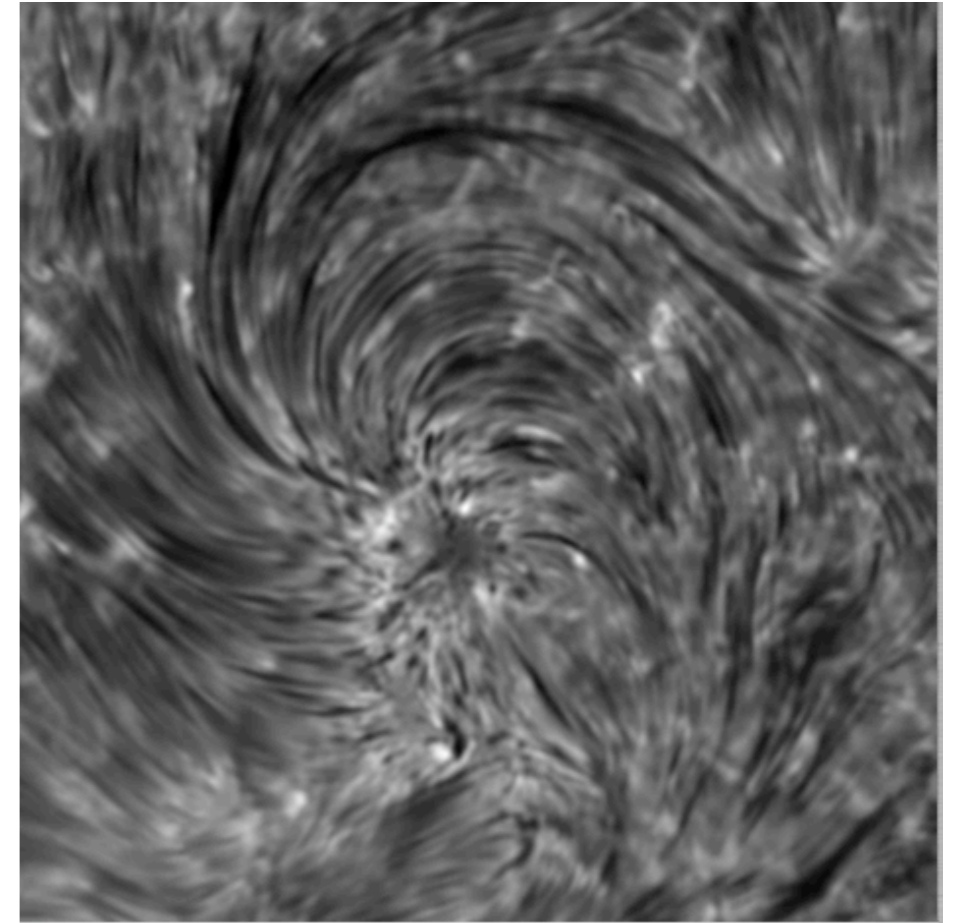
Inspiration for  
much work  
generically  
called  
“chromospheric  
fine structure”

**magnetic interface**  
**physical considerations**

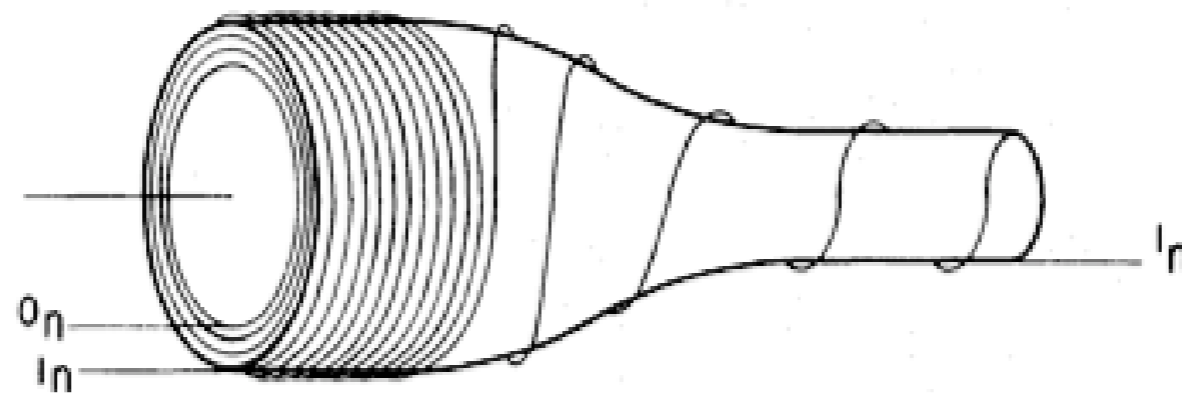


# Note: twist/ electrical currents can be easier to detect in the chromosphere!

- IBIS again: *clear  $B_\phi \Rightarrow j_z$*
- Hinode rotating spicules
- ang. mom. conservation around tubes
- Knölker et. al. (1988)- tube stability requires rotating flow
- Parker (1974):  *$B_\phi/B_z$  increases with  $z$*



## DYNAMICAL PROPERTIES OF MAGNETIC FIELD



# Chromosphere vs. photosphere as the coronal boundary

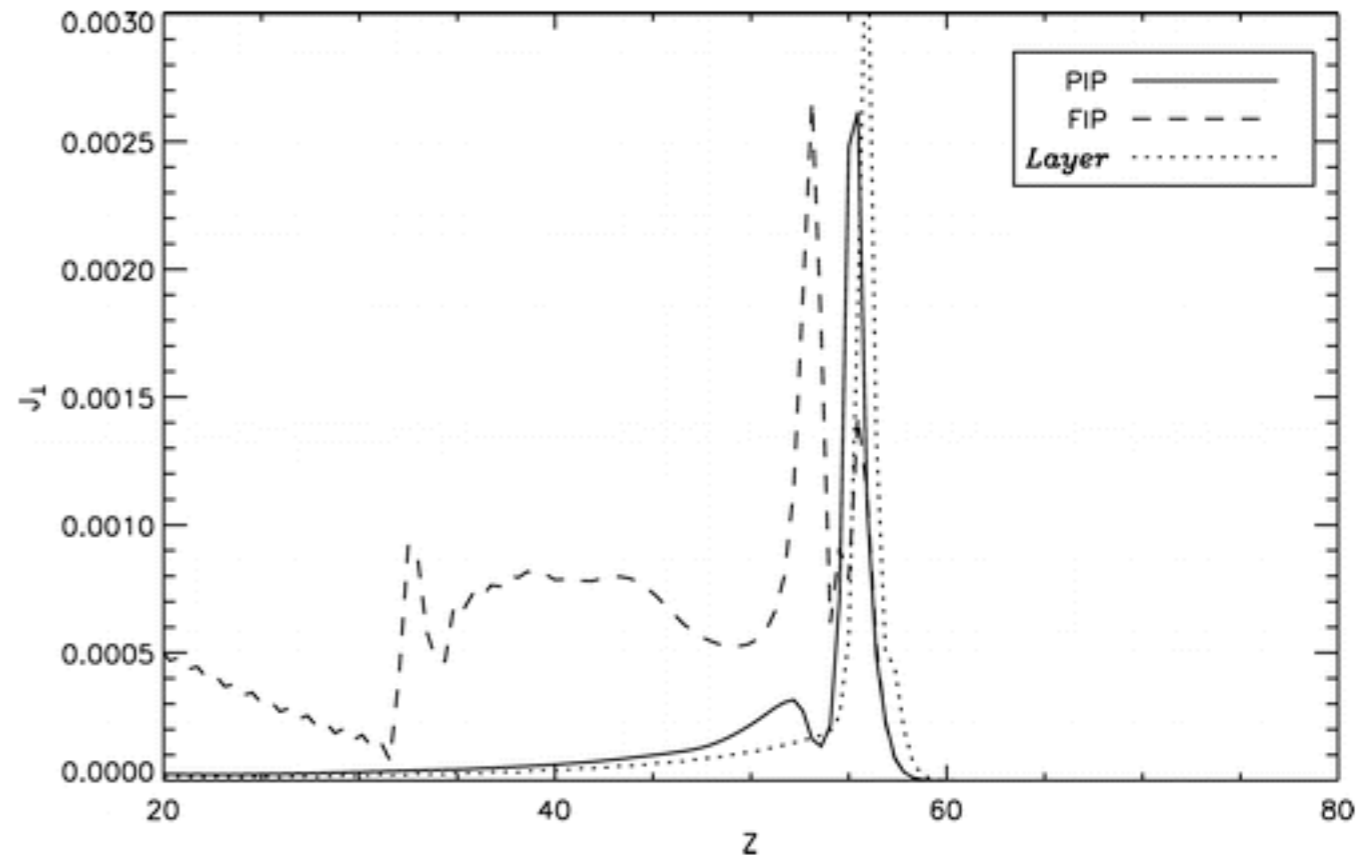
- chromosphere spans 9 scale heights
- $\Rightarrow$  chromosphere usually contains  $\beta=1$  surface
  - $|\mathbf{j} \times \mathbf{B}| \rightarrow \beta B^2/2\mu$  above  $\beta=1$      $\mathbf{j}_\perp \rightarrow$  **small**
- partial ioniz<sup>n</sup>  $\Rightarrow$  3-fluid frictional dissipation, **heating**
  - $Q_{\text{fr}} = \mathbf{j}^2/\sigma + (\xi_n \mathbf{j} \times \mathbf{B} - \mathbf{G})^2/\alpha_n$ ,     $\mathbf{G} = \xi_n \nabla p - \nabla p_n$
  - “ambipolar diffusion”/star formation (1950s [Schlüter, Cowling](#))
- case  $\mathbf{G} = \mathbf{0} \Rightarrow$  “Cowling conductivity”  $\sigma_\perp^*$  ([Arber & cohorts](#))
  - $Q_{\text{fr}} = j_{\parallel}^2/\sigma + j_\perp^2/\sigma_\perp^*$      $\sigma/\sigma_\perp^* = 1 + 2 \xi_n \omega_e \tau_e \omega_i \tau_i \gg 1$
  - $\Rightarrow$  **dissipation of  $\mathbf{j}_\perp$** .    *Explains why IBIS nearly f-f?*
- **NOTE:**  $\sigma_\perp^*$  is some steps removed from  $\sigma$  (kinetic theory)
  - case  $\mathbf{G} \neq \mathbf{0}$ :  $\sigma_\perp^*$  incorrect!
  - one must simultaneously determine the nature of  $\mathbf{j}_\perp$  (cf. E-region electrojet) from the dynamics

# Chromosphere tends to “filter out” $j_{\perp}$ : coronal base magnetic field $\rightarrow$ force-free

- Braginskii (1965): certain motions ( $\mathbf{G}...$ ) dissipate  $j_{\perp}$ 
    - Alfvén, fast modes, dynamic situations where  $\nabla p - \rho \mathbf{g} + \mathbf{j} \times \mathbf{B} \neq \mathbf{0}$
  - **Not** slow modes, slow dynamics (cf. Goodman 2000)
  - So, at coronal lower boundary, chromosphere makes:
    - $j_{\perp} \sim 0$ ;  $\mathbf{j} \times \mathbf{B} \sim 0$
    - weaker Alfvén/fast modes
    - $\text{curl } \mathbf{B} = \alpha \mathbf{B}$ :  $\alpha(r) \rightarrow \text{constant?}$
- (Parker current sheets..)

*Flux emergence*: Arber, Haynes & Leake (2007) based upon Cowling’s conductivity ( $\mathbf{G}=\mathbf{0}$ ):

Plot of the magnitude of  $j_{\perp}$  as a function of height along the line  $x = y = 0$  for all three resistivity models at  $t = 160$ .

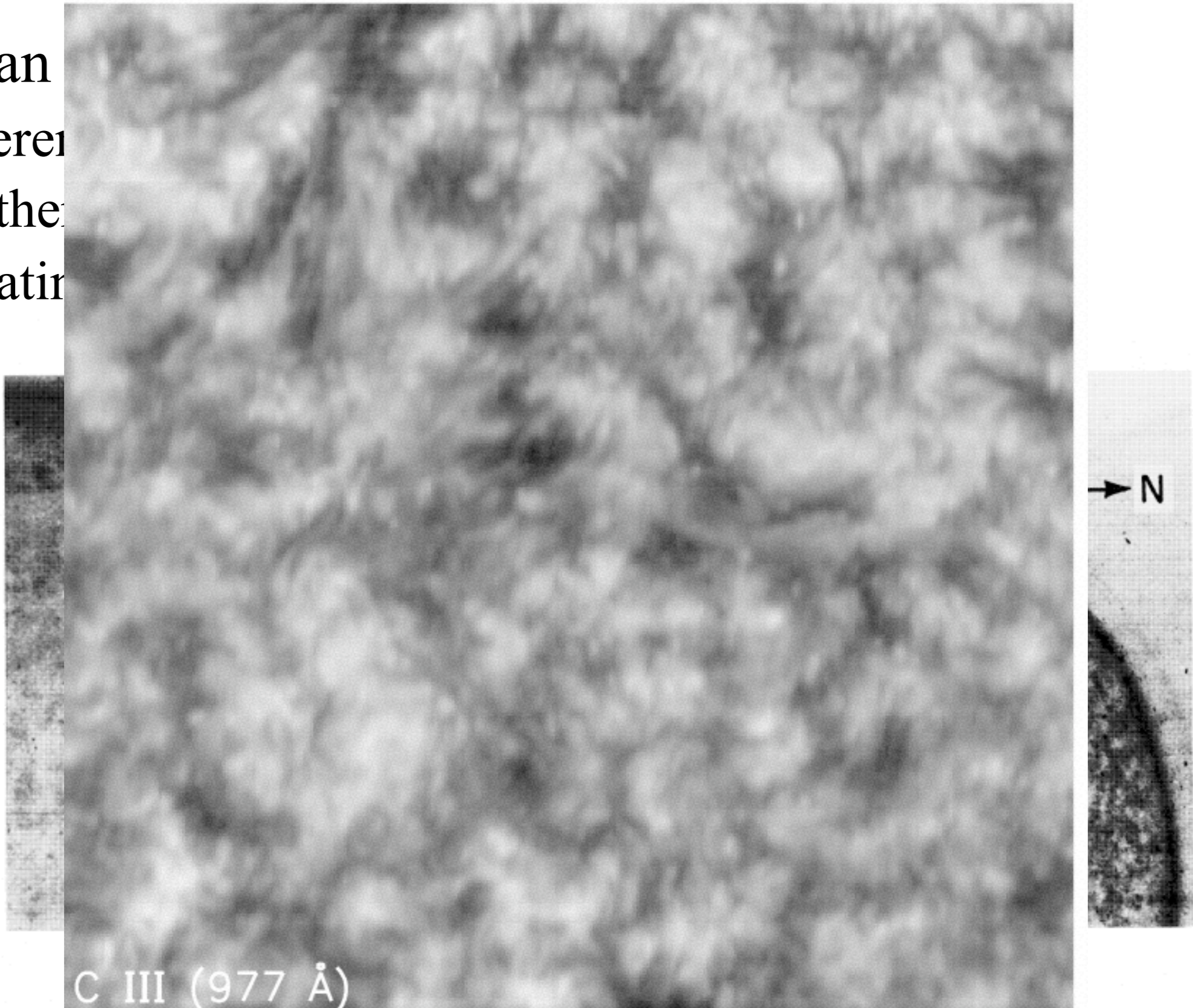


...radical effect on flux emergence process

# **thermal interface**

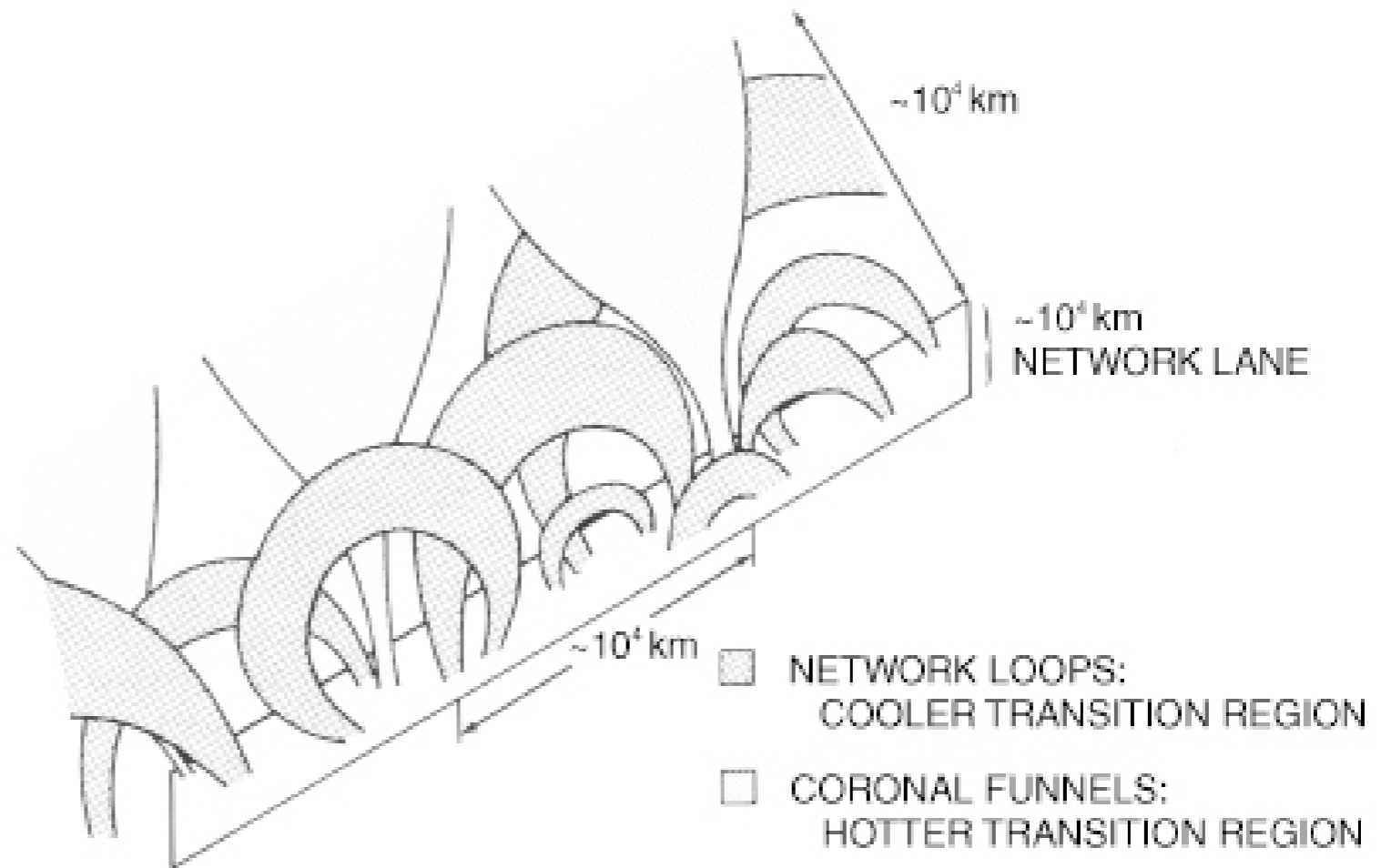
# The problem- observations

- Feldman
  - differen
  - TR the
  - radiati



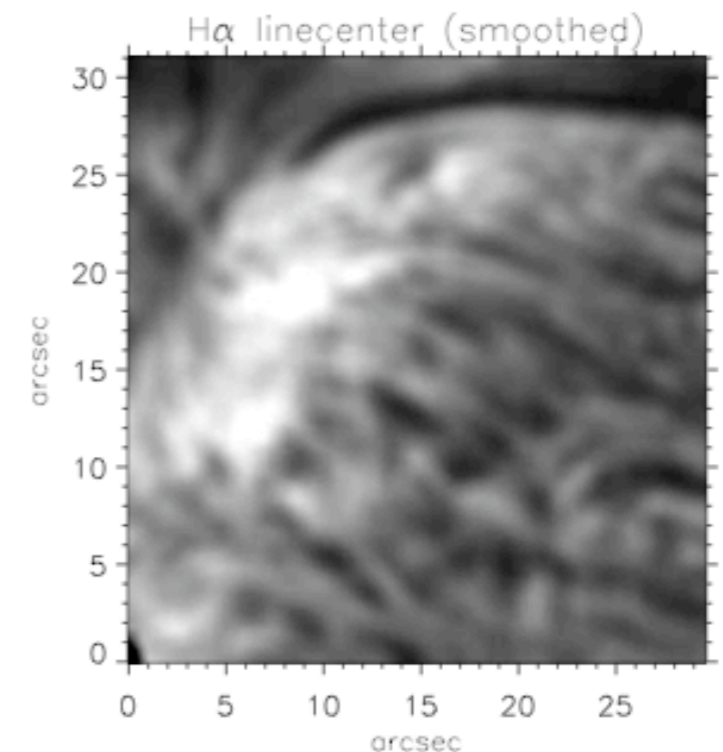
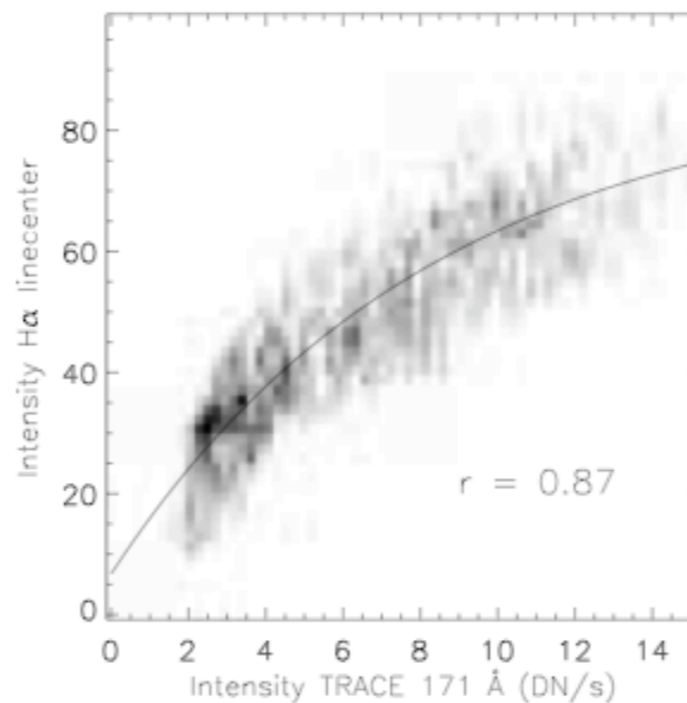
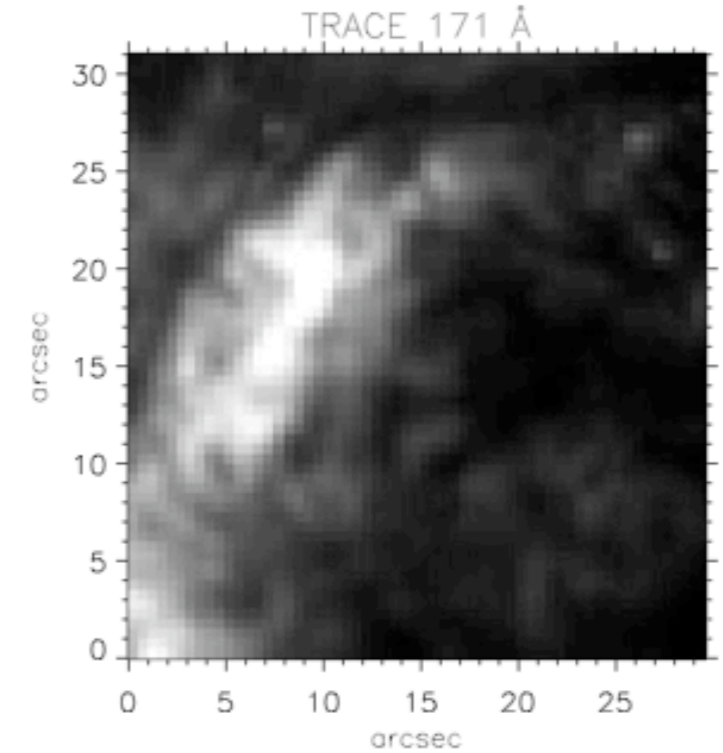
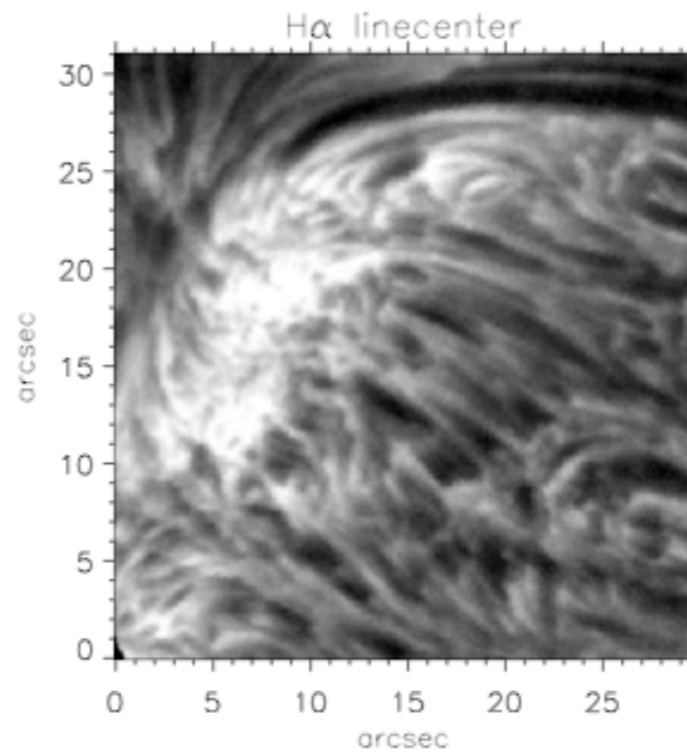
## Dowdy et al. (1986)

- Mixed polarity within network boundaries
- tries to explain “UFS”
- indeed these are thermally and magnetically separate entities



# Depontieu et al 2003: TRACE/SST data

## CORRELATIONS BETWEEN CHROMOSPHERIC AND TR EMISSION



Yet...

Significant correlations exist between the H $\alpha$  chromospheric intensity and the low corona

# Questions concerning cool loops

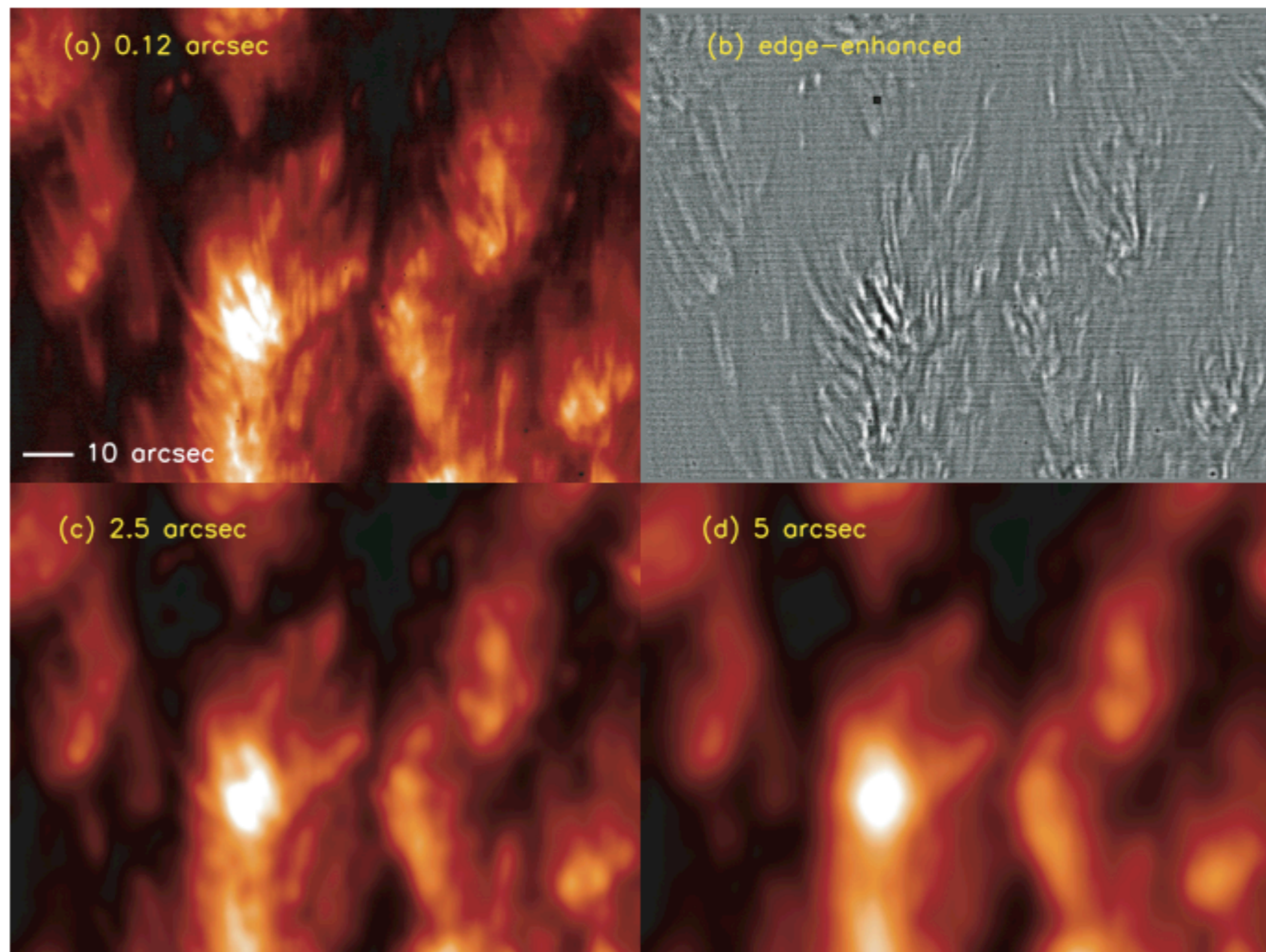
- Cool loops are considered by most a viable explanation, but
- where does the  $10^6 \text{ erg cm}^{-2} \text{ s}^{-1}$  conductive flux go?
- Is it merely a coincidence that the lower TR radiates about  $10^6 \text{ erg cm}^{-2} \text{ s}^{-1}$ ?
- Why should the cool loop distribution make the upper (conductive) and lower (cool loop) TR be correlated, at least on scales  $>$  a few Mm?
- are they stable (Cally & Robb 1991)?
- where are the tell-tale magnetic footpoints?
- ...



# Judge & Centeno (2008)

- VAULT  $L\alpha$  data vs. KPNO magnetic data
  - supplemented by Hinode SP vector polarimetry
- Prompted by Patsourakos et al (2007)
  - We noted something “odd” about proposed cool loops
  - **large-scale alignment of  $L\alpha$  threads**

Patsourakos et al:



# KPVT+POTL FIELDS+VAULT

## active network

Black=low-lying loops ( $h < 5$  Mm)

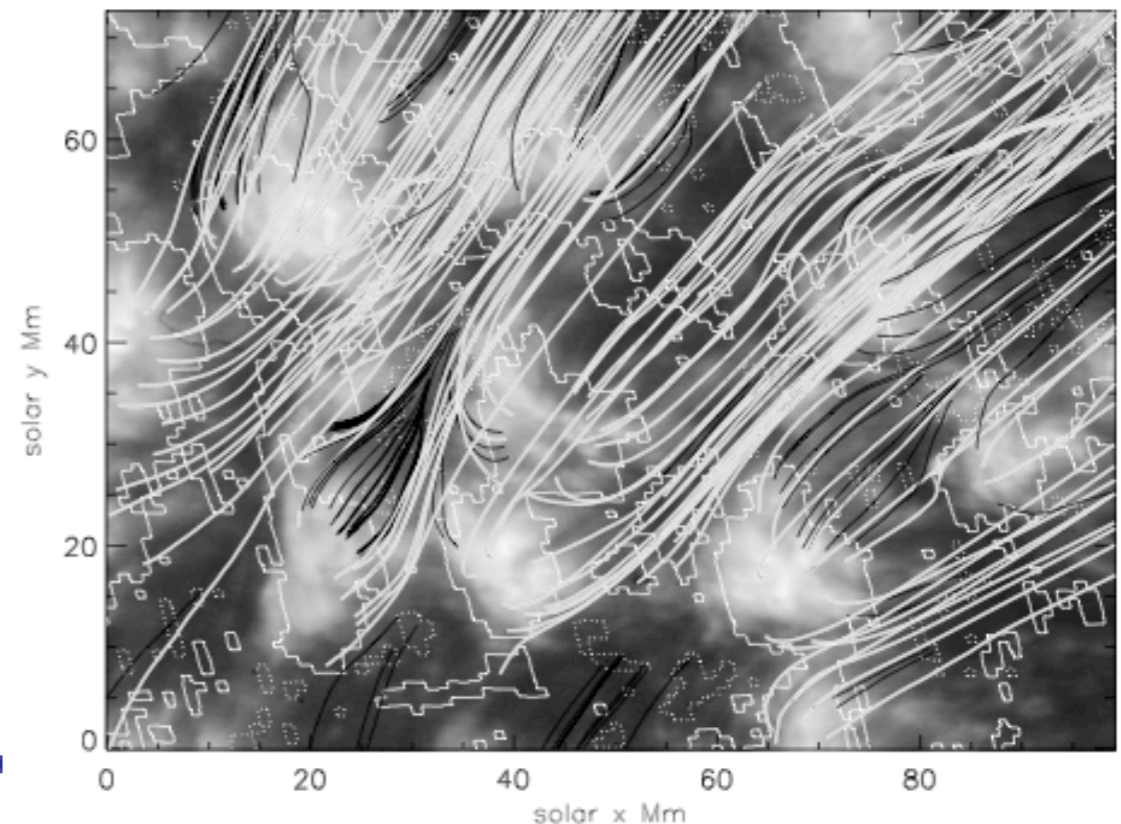
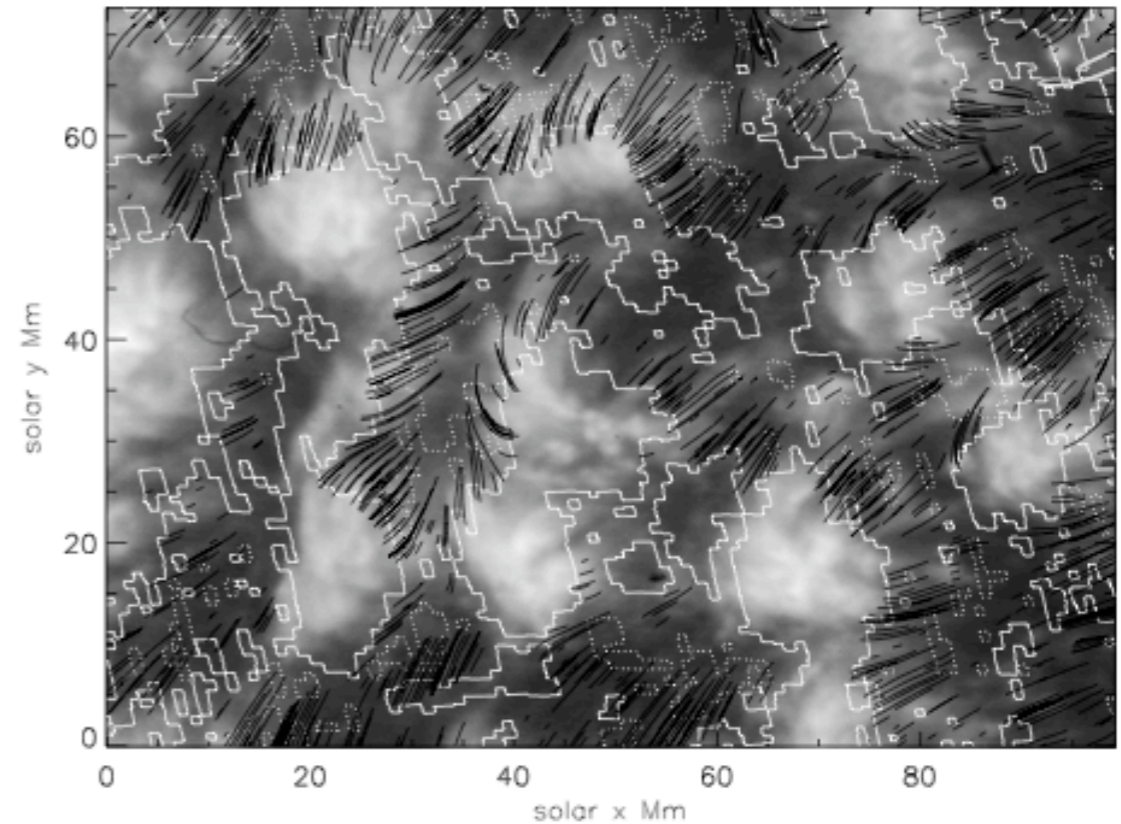
Gray= long

Stability requires that low-lying loops are possibly cool, but higher loops must be hot

Most  $L\alpha$  emission originates from the base of hot, coronal loops

Some may arise from cool loops, but not commonly in active network

Cannot appeal to “unresolved (salt +pepper) fields”-  $L\alpha$  emission forms above  $h=0.8$  Mm. “Loops” with footpoints separated by 1” can’t reach these heights

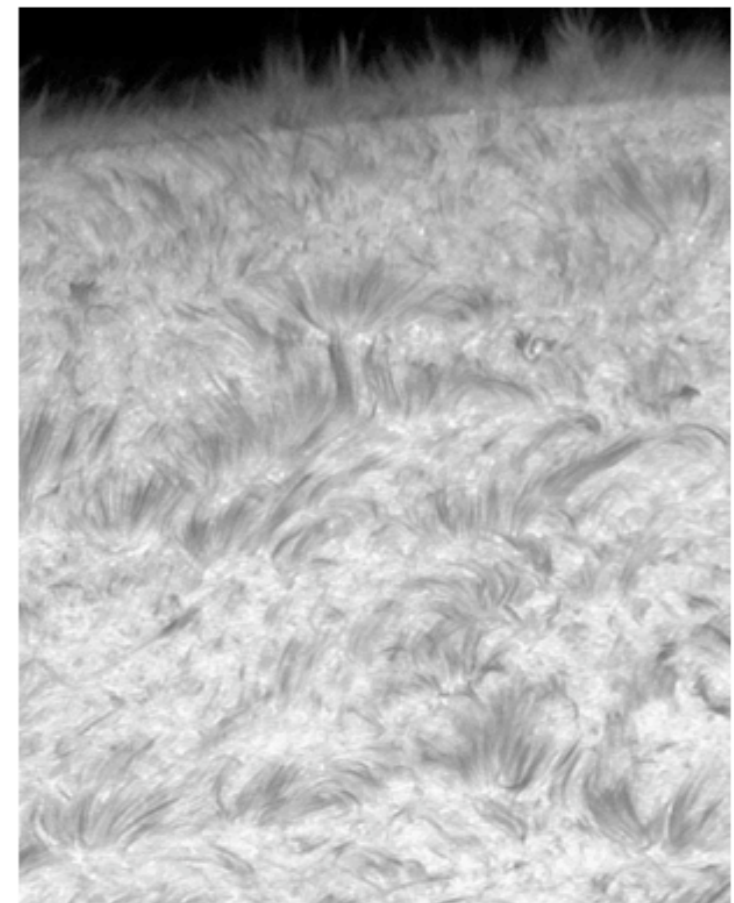
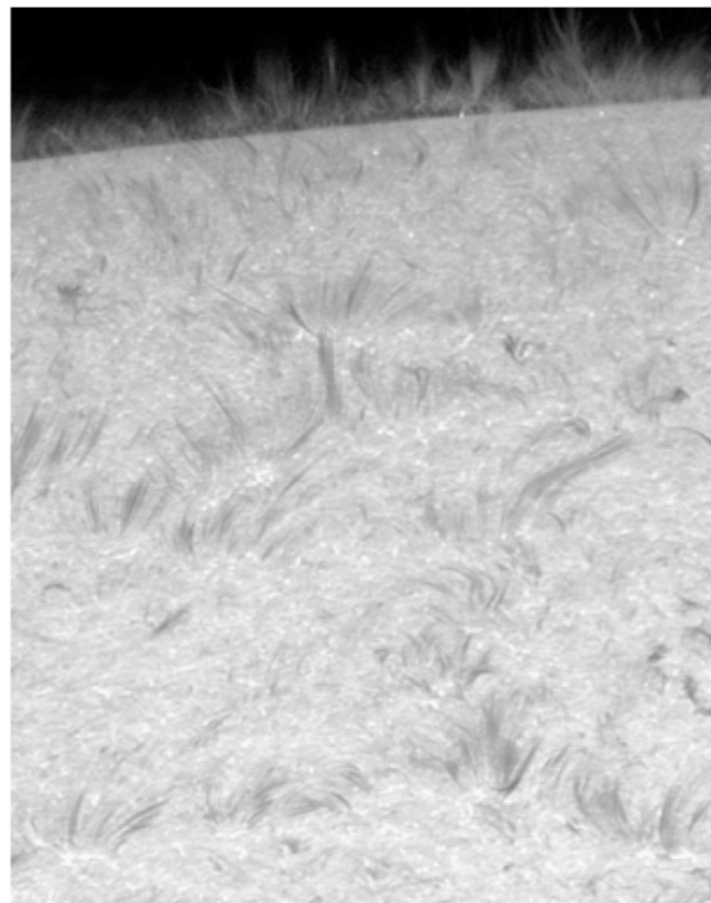
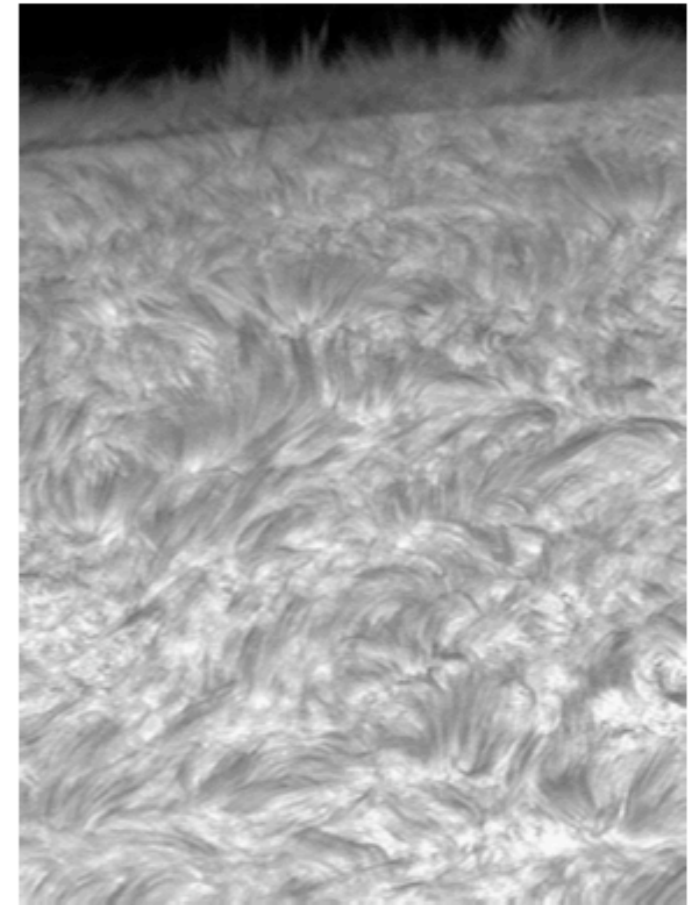
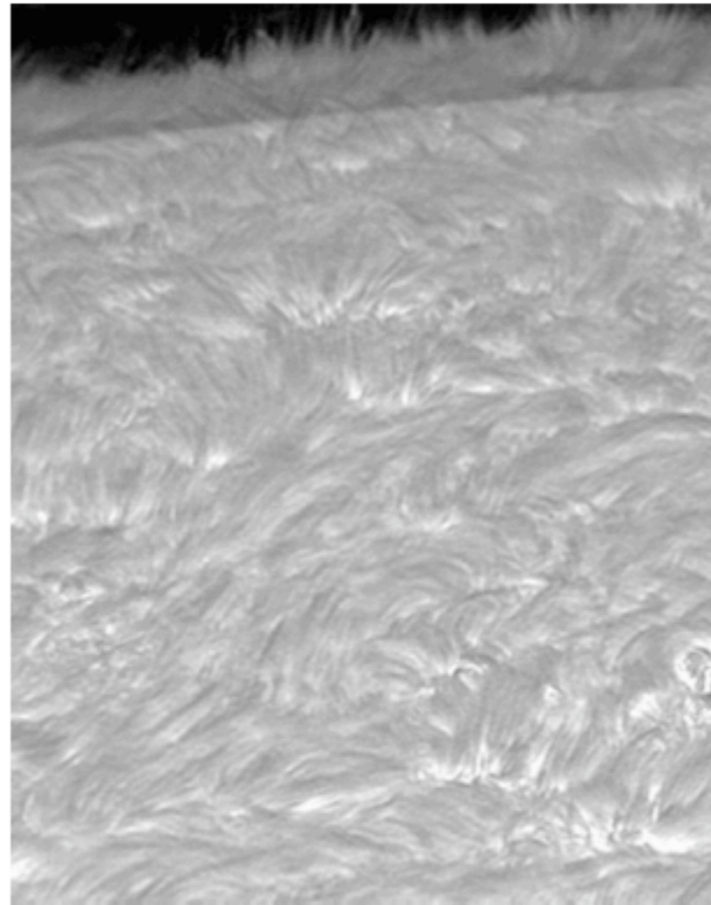


# Spicules, fibrils..

- base of the corona is a **non-planar** thermal boundary
- e.g., DOT H $\alpha$  (Rutten 2007) clockwise 0, -0.4, -0.6, -0.8 Å:

*consider  $\alpha$  in*

*$\text{curl } \mathbf{B} = \alpha \mathbf{B}$  for photosphere  
and coronal base*

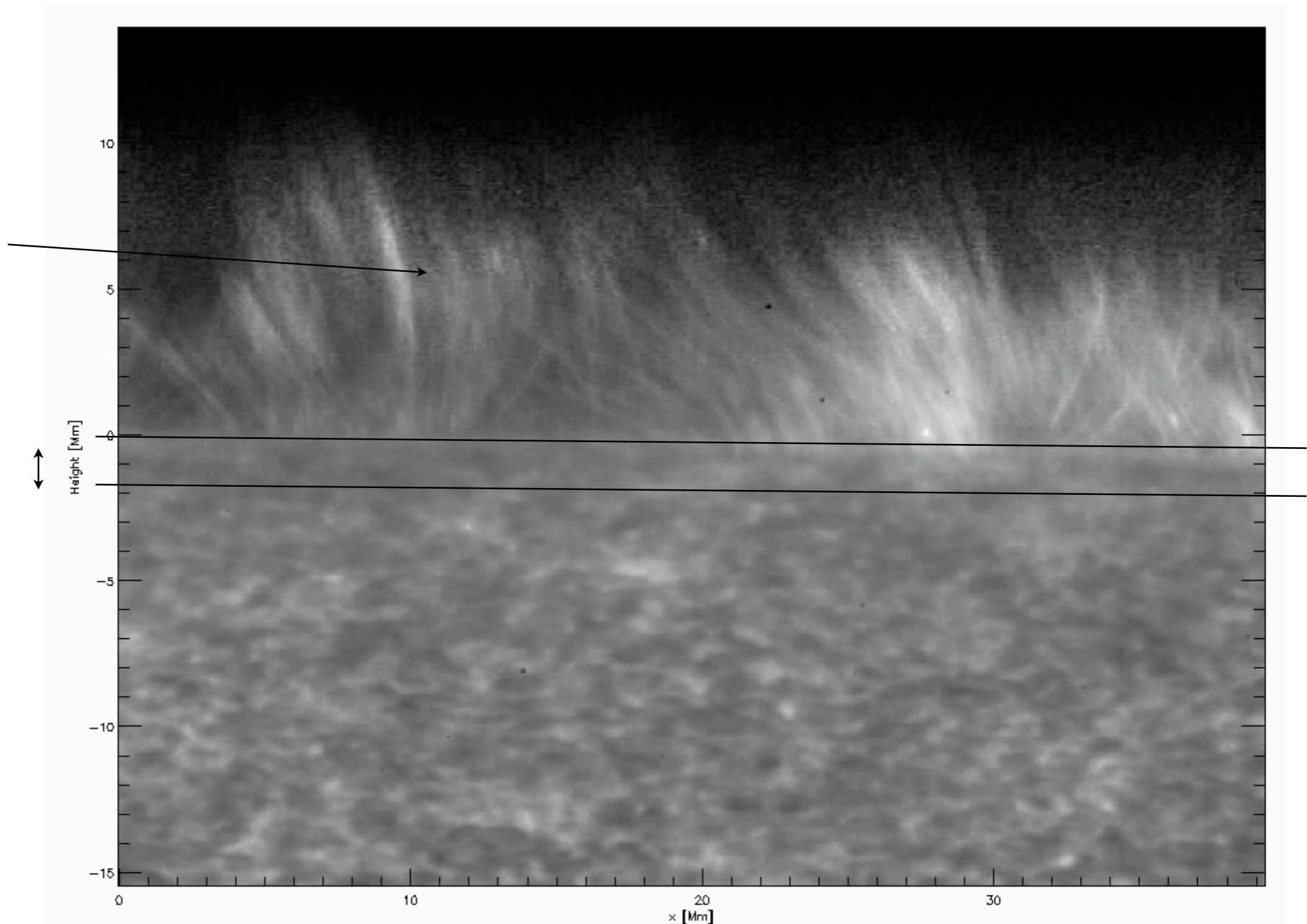


# Hinode spicules

- Ca II (radial filter to enhance spicules, M. Carlsson)

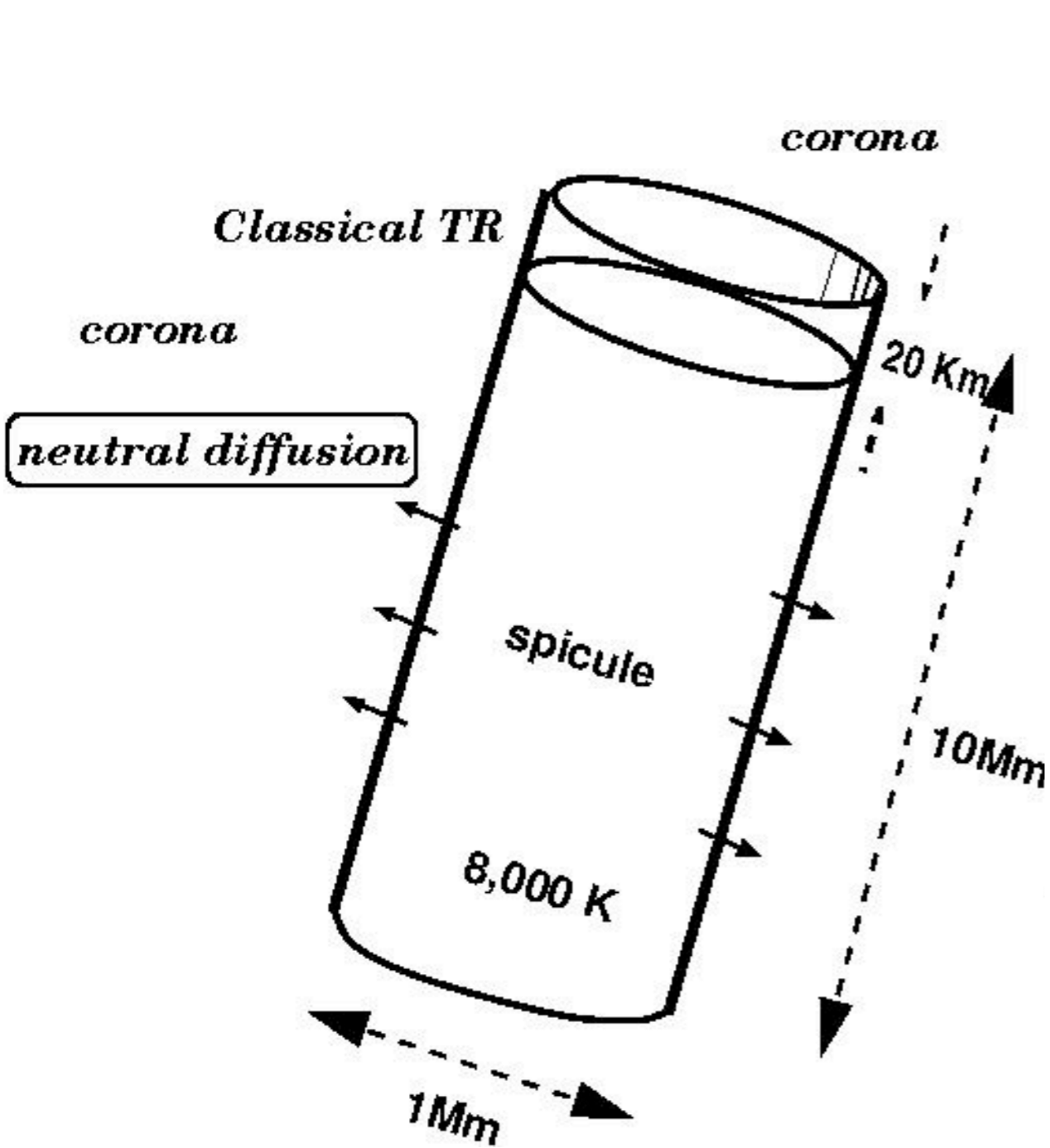
spicules *arise*  
*from within*  
the chromo-  
sphere

stratified VAL  
chromosphere  
1.5Mm only



# Judge (2008) ApJL 683, 87-90

## “spicule” → cross field diffusion → TR radiation



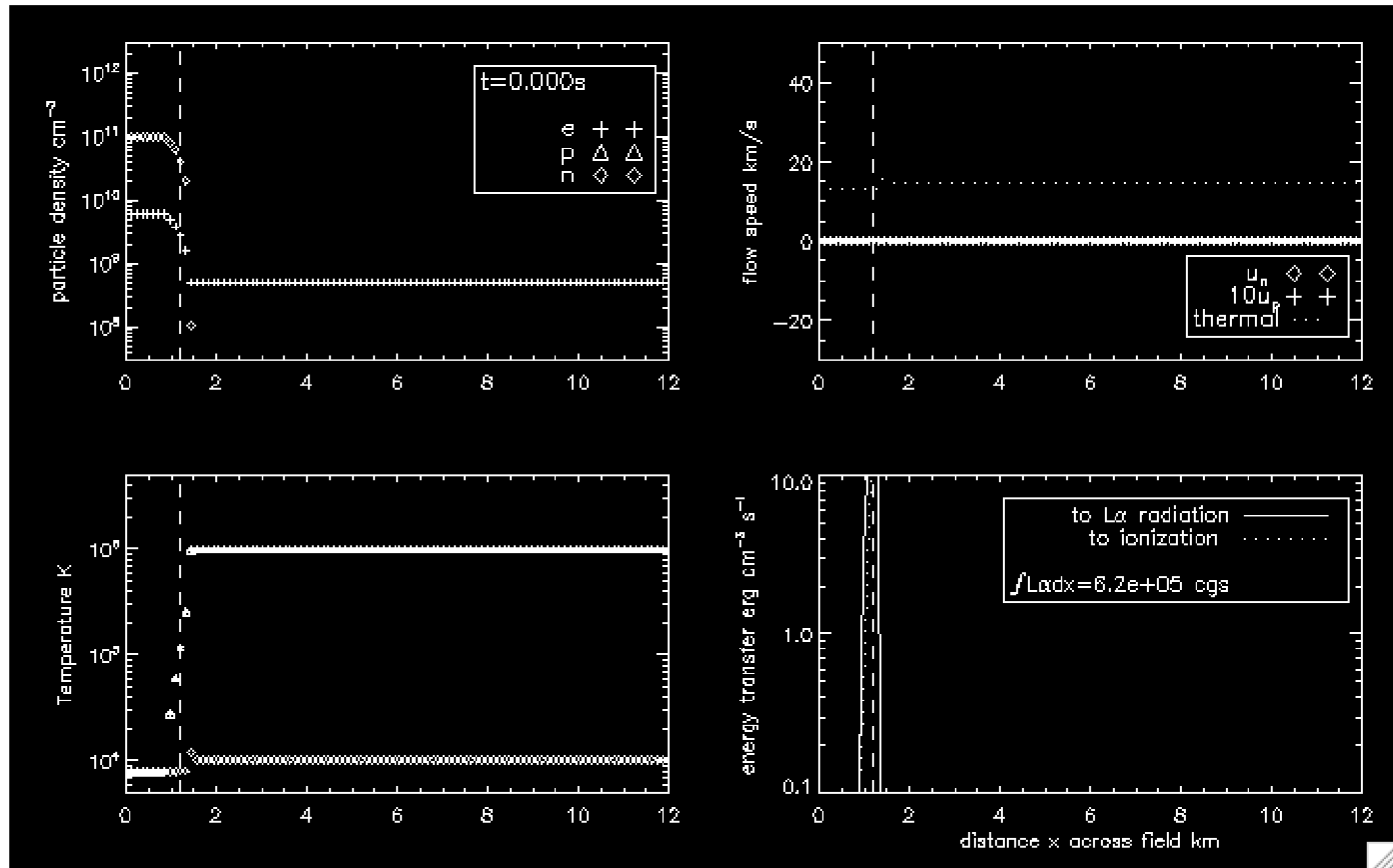
Initial corona

$T_h$	K	$10^6$	
$n_h$	$\text{cm}^{-3}$	$8.0 \times 10^8$	
$n_p, n_e$	$\text{cm}^{-3}$	$4.0 \times 10^8$	
$p$	$\text{cm}^{-3}$	$1.1 \times 10^{-1}$	
$B$	G	10	
$\beta$		$2.8 \times 10^{-2}$	
$\omega_p$	$\text{s}^{-1}$	$9.6 \times 10^4$	
$\tau_{gyro}$	km	$1.5 \times 10^{-3}$	
$\tau_{pp}$	s	1.6	$n_p^{-1} T^{+3/2}$
$\omega_p \tau_{pp}$		$1.5 \times 10^5$	
$\tau_{ee}$	s	$5.0 \times 10^{-2}$	$n_e^{-1} T^{3/2}$
chromospheric tube			
$T_c$	K	$8.0 \times 10^3$	
$\bar{v}$	$\text{km s}^{-1}$	13	$T^{1/2}$
$n_c$	$\text{cm}^{-3}$	$10^{11}$	
$\tau_{nn}$	s	$1.4 \times 10^{-2}$	$n_n^{-1} T^{-1/2}$

# Results: model $L\alpha \sim 0.1x$ observed using only local coronal heat

1D 3-fluid  
calculation  
of cross-field  
diffusion  
from a cool  
flux tube into  
coronal  
plasma

no field  
aligned  
conduction



calculations with different coronal  $n, T$ : non-linear  
relationship between  $L\alpha$  and coronal emission

# Judge (2008)

- calculations for  $L\alpha$  are promising, (also  $L\beta$ , He I 584)
  - this is the hardest line to explain, others may follow?
- **cross-field diffusion of neutrals** might solve the 40+ yr problem of **energy balance in extended structures in the lower TR**
- chromosphere supplies the mass, corona the energy
  - cool loops don't explain active network (Judge & Centeno 2008)
  - “UFS” in this new picture is **thermally connected to the corona**
- needed
  - 2D calculations including field-aligned conduction and dynamics
  - observations of the chromosphere/corona interface in relation to magnetic field

# Conclusions

- the magnetic chromosphere remains poorly understood
- the Sun undergoes the awkward transition from forced  $\beta > 1$  to force-free  $\beta < 1$  there:  $\mathbf{j} \times \mathbf{B} \rightarrow \mathbf{0}$  at the coronal base
- magnetic free energy  $\rightarrow$  chromospheric heat and radiation
  - dissipation of  $\mathbf{j}_\perp$ :  $\mathbf{j} \times \mathbf{B} \rightarrow \mathbf{0}$ ,
  - $\alpha(\mathbf{r}) \rightarrow ?$  at the coronal base: Parker's current sheets
  - observed chromospheric losses might arise from  $\mathbf{j}_\perp \cdot \mathbf{E}$ ? (friction)
- spicules/fibrils+neutral diffusion+coronal heat finally explains the transition region?
- **meaningful** photos./chromos. polarimetry is here and is needed to
  - understand basic MHD physics (e.g. Pietarila & colleagues)
  - understand magnetism at the coronal base (e.g. Wiegelmann, Schrijver)
- 3-fluid MHD models are needed to assess how chromospheric processes influence the coronal base conditions (e.g., validity of Cowling's  $\sigma_\perp^*$ )



# To understand the corona we must understand what is under Gold's line... *is single-fluid MHD adequate?*

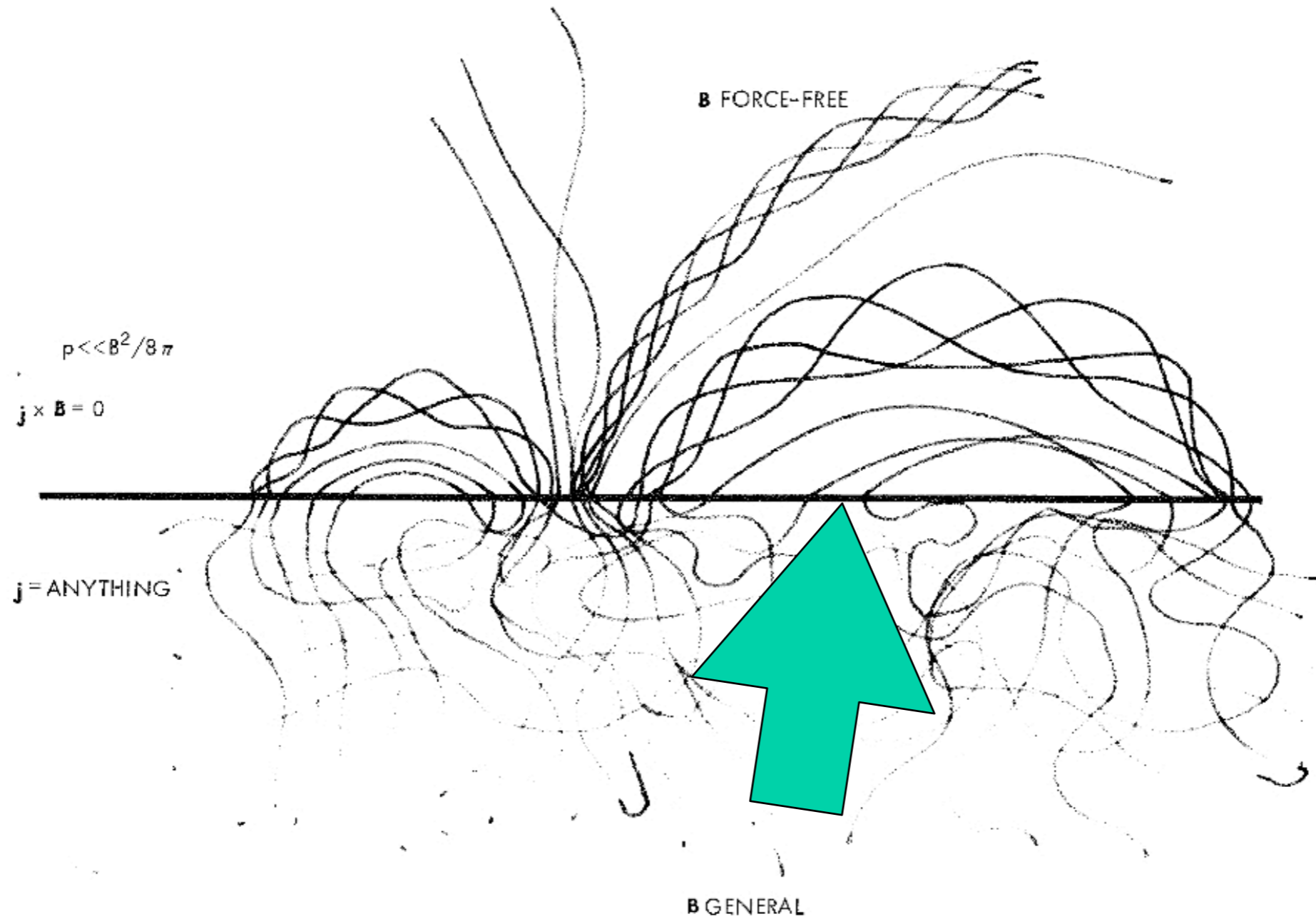


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